



## Abstract

**Purpose:** The alliance of charged particle therapy and the spatial fractionation of the dose, as in minibeam or Grid therapy, is an innovative strategy to improve the therapeutic index for the treatment of radioresistant tumors. The aim of this work was to assess the optimum irradiation configuration in heavy ion spatially fractionated radiotherapy (SFRT) in terms of ion species, beam width, center-to-center distances, and linear energy transfer (LET), information that could be used to guide the design of the future biological experiments. The nuclear fragmentation leading to peak and valley regions composed of different secondary particles, creates the need for a more complete dosimetric description than the classical one in SFRT.

**Methods:** Monte Carlo simulations (GATE 6.2) were performed to evaluate the dose distributions for different ions, beam widths and spacings. We have also assessed the 3D-maps of dose-averaged LET and proposed a new parameter, the peak-to-valley-LET ratio, to offer a more thorough physical evaluation of the technique.

**Results:** Our results show that beam widths larger than 400  $\mu\text{m}$  are needed in order to keep a ratio between the dose in the entrance and the dose in the target of the same order as in conventional irradiations. A large ctc distance (3500  $\mu\text{m}$ ) would favor tissue sparing since it provides higher PVDR, it leads to a reduced contribution of the heavier nuclear fragments and a LET value in the valleys a factor 2 lower than the LET in the ctc leading to homogeneous distributions in the target.

**Conclusions:** Heavy ions MBRT provide advantageous dose distributions. Thanks to the reduced lateral scattering, the use of submillimetric beams still allows to keep a ratio between the dose in the entrance and the dose in the target of the same order as in conventional irradiations. Large ctc distances (3500  $\mu\text{m}$ ) should be preferred since they lead to valley doses composed of lighter

nuclear fragments resulting in a much reduced dose-averaged LET values in normal tissue, favoring its preservation. Among the different ions species evaluated, Ne stands out as the one leading to the best balance between high PVDR and PVL<sub>R</sub> in normal tissues and high LET values (close to 100 keV/ $\mu$ m) and a favourable oxygen enhancement ratio in the target region.

7 Keywords: Hadron Minibeam radiation therapy, Grid therapy, Monte Carlo simulations

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## 8 I. INTRODUCTION

9 The use of the spatial fractionation of the dose in charged particle therapy is a new  
10 therapeutic approach [1], [2] and [3] to overcome the main limitation in radiotherapy (RT),  
11 namely the normal tissue tolerances. Spatial fractionated radiotherapy (SFRT) started soon  
12 after the birth of RT with the aim of reducing skin damage when deep seated tumors were to  
13 be treated with the existing orthovoltage machines. This approach, called GRID therapy [4],  
14 is currently used clinically at few hospitals using megavoltage (MV) photon beams delivered  
15 by medical Linacs [5], [6] and [7]. GRID therapy (GRT) has been successfully used to deposit  
16 high doses in the treatment of bulky tumors, for which conventional approaches fail in their  
17 management [6], [7] and [8]. However, the important lateral scattering of MV beams and  
18 the relative low fluence of Linacs limits the gain of this technique [9]. The extremely high  
19 fluence of non-divergent kilovoltage beams provided by third generation synchrotron sources  
20 triggered the alliance of the spatial fractionation of the dose with the use of very narrow  
21 (submillimetric) beams, as in microbeam and minibeam radiation therapies [10], [11] and  
22 [12]. The use of small field sizes allows to better exploit the dose-volume effects [13], further  
23 increasing the normal tissue sparing. Indeed, extremely high normal tissues dose tolerances  
24 have been observed in animal studies [10], [14], [15], [11] and [16]. Tumor growth delay  
25 in aggressive animal tumor models was also observed [17], [18], [19], [20] and [21]. These  
26 outcomes seem to be the result of the participation of some biological mechanisms (not well  
27 understood yet) different from those ones in ‘standard’ RT.

28 A further improvement of therapeutic index could be achieved by joining the advantages  
29 of the spatial fractionation of the dose to the high dose conformity and increased relative  
30 biological effectiveness (RBE) of heavy ions. The exploration of this new approach has re-

31 cently started. Different particles have been studied. Proton MBRT [2] has already been  
32 implemented at a clinical center [22] and the first biological experiments confirms a remark-  
33 able gain in normal tissue tolerances [23] and [24] reported the evaluation (Monte Carlo  
34 simulations) of the broadening in water of several light ions (from protons to  ${}^7\text{Li}$ ). Among  
35 the different types of ion species, carbon ions are considered to have optimal properties in  
36 terms of superior physical and biological characteristics [25]. The results of the recent stud-  
37 ies on carbon MBRT from dosimetric [26], [27] and biological point of view [1], encourages a  
38 further exploration of this approach. Due to their increased linear energy transfer, heavier  
39 ions, like oxygen, are considered specially promising for the treatment of hypoxic tumors,  
40 which remains one of the major challenges in radiotherapy (RT). Radiobiological findings  
41 in the laboratory indicated that resistant cells of hypoxic tumors could be effectively de-  
42 stroyed with very heavy ion beams, such as silicon and argon [28]. However, clinical results  
43 with a few patients performed with argon in 1979 and with silicon in 1982 lead to adverse  
44 late tissue results, and the use of these beams was discontinued. Nevertheless, the gain in  
45 tissue sparing that might be provided by the use minibeam radiation therapy might allow  
46 profiting from the remarkable effectiveness of very heavy ions. This could offer a new hope  
47 for aggressive hypoxic tumors, whose treatment with conventional methods is very limited.  
48 The advantageous dose distributions obtained in the very first Monte Carlo (MC) study [3]  
49 support the further exploration of this avenue.

50 This work is devoted to the assessment of optimum irradiation configurations in heavy  
51 ion spatially fractionated RT (SFRT) in terms of ion species, beam width and shape, and  
52 center-to-center (ctc) distance, that could be used to guide the design of future biological  
53 experiments. In order to maximize the dose volume effects, we have evaluated for each  
54 ion the minimum beam width that could be used without significantly degrading the ratio

55 between the entrance dose and the dose at the Bragg peak [29].

56 In SFRT, the resulting dose profiles follow a pattern of peaks and valleys, with high doses  
57 along the primary beam path (peaks) and low doses in the space between them (valleys).  
58 Previous experiments at synchrotrons (x-rays) indicate that normal tissue sparing seems to  
59 benefit from a high peak-to-valley dose ratio (PVDR) and low valleys while low PVDR and  
60 high doses to the tumor favor tumor control [18]. PVDR is used as an important dosimetric  
61 variable in x rays SFRT. However, due to nuclear fragmentation, PVDR or valley dose,  
62 may not provide a thorough description in heavy ion SFRT. Indeed, the produced nuclear  
63 fragments are going to fill the valleys, which will then be composed of a mix of particles.  
64 The same PVDR or even the same valley dose can correspond to different composition of  
65 fragments leading to different effects on normal tissues. This is usually taken in considera-  
66 tion in heavy ion therapy by using RBEs. However, uncertainty on the RBE is often quoted  
67 as a major hindrance to a widespread use of heavy ions in radiotherapy and is a source  
68 of concern for the potential late effects. For instance, C-ions have been used very little in  
69 pediatric patients mostly for the concern about high risk of secondary cancers. In addition,  
70 the expected radiobiological features of SFRT are so different from those of ‘standard’ irra-  
71 diations, probably including some non-targeted effects that the radiobiological effect cannot  
72 be assessed by means of established RBEs of ‘standard’ hadron therapy (seamless beam  
73 irradiation). Biological experiments are needed to assess the effectiveness of this approach  
74 and reconstruct the complete picture. In the meantime, maps of dose-averaged Linear En-  
75 ergy Transfer (LET) can be reliably quantified by means of MC simulations [30] and used to  
76 predict the regions of expected higher biological effect. LET describes the density of ioniza-  
77 tion and, therefore, the probability of the damage. New parameters based on dose-averaged  
78 LET, more suitable for a quantitative physical description of heavy ion therapy SFRT, have

79 been defined and evaluated.

80 Since our approach is especially suitable for tumors for which pulmonary and/or cardiac  
81 cycles (target motions) have minimal effects, in particular, for presently difficult-to-treat  
82 neurological indications, our study was focused on parameters adequate for brain tumor  
83 treatments.

## 84 II. MATERIALS AND METHODS

### 85 A. Monte Carlo simulations

86 The code GATE (version 6.2) [31] was employed in this work. GATE is an open source  
87 MC simulation platform based on GEANT4 (version 4.9) [32] enabling the modeling of emis-  
88 sion and transmission tomography, and radiotherapy. The option 3 of the electromagnetic  
89 standard physics package for GEANT4 and the recommended hadronic physics list for carbon  
90 were used [32]. A range cut value of  $50 \mu\text{m}$  was used for all particles in the water phantom.  
91 The number of primary particles employed was  $4 \times 10^9$  ions. This leads to 1% of maximum  
92 uncertainty (with a cover factor  $k = 2$ ).

93 A magnetic collimation was assumed to generate minibeam of six different ions ( $^{12}\text{C}$ ,  $^{16}\text{O}$ ,  
94  $^{20}\text{Ne}$ ,  $^{28}\text{Si}$ ,  $^{40}\text{Ar}$  and  $^{56}\text{Fe}$ ). Two different irradiation geometries were considered: rectangular  
95 (1D-Grid or minibeam (MBRT)) and squared (2D-Grid (GRT)) beams, with dimensions  
96 ranging from  $100 \mu\text{m} \times 2 \text{ cm}$  to  $700 \mu\text{m} \times 2 \text{ cm}$  in the first case and from  $300 \mu\text{m} \times 300 \mu\text{m}$   
97 to  $1000 \mu\text{m} \times 1000 \mu\text{m}$ , in the second case. See figure 1. The sources were located at  
98 the phantom entrance. The beams impinged on a cylindrical water phantom whose size  
99 (height and diameter of 16 cm) mimics the human head and it has been used in previous  
100 dosimetric studies [3] and [27]. A virtual tumor ( $2 \text{ cm} \times 2 \text{ cm} \times 2 \text{ cm}$ ) was placed at 7

101 cm-depth. A realistic angular spread of 3 mrad [33] and [34] and Gaussian shape were  
 102 considered. The dose accumulation voxels had dimensions of  $60 \mu\text{m} \times 2 \text{mm} \times 1\text{mm}$  and  
 103  $60 \mu\text{m} \times 60 \mu\text{m} \times 1\text{mm}$  for MBRT and GRT, respectively. To save computation time only  
 104 one single beam was simulated and the dose distributions were evaluated as the sum of the  
 105 contributions of each individual minibeam to cover the virtual tumor ( $2 \text{cm} \times 2 \text{cm}$ ).

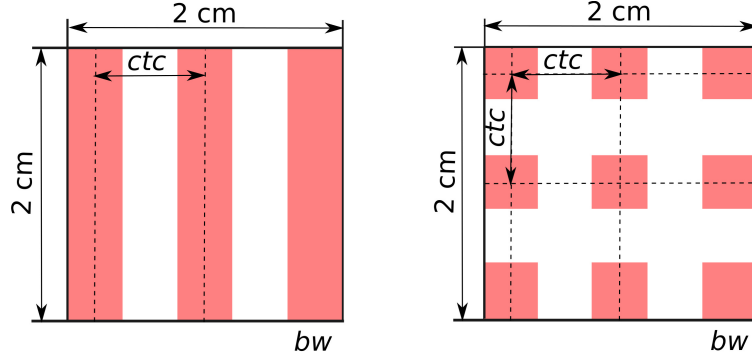


FIG. 1: Scheme representing the geometrical parameters in MBRT (left) and in GRT (right). *ctc* refers to the distance between two consecutive beam centers and *bw* to the beam width.

106 First, pristine Bragg peak curves were calculated to determine the minimum beam width  
 107 (*bw*) for each ion/configuration by evaluating the ratio between the dose at the entrance  
 108 and the dose at the Bragg Peak. A maximum ratio of around 25 %, similar to the ones that  
 109 would be obtained in conventional (broad beam) irradiations [35], was used as a criterion.  
 110 Then, SOBP irradiations covering the target homogeneously in depth were also simulated.  
 111 With that aim 50 monoenergetic ion beams were weighted using the method described by  
 112 Jette and Chen [36]. The table II A show the minimum and maximum energies used for  
 113 each ion.

114 The peak-to-valley dose ratio as a function of depth was assessed for each ion and its  
 115 corresponding minimum width. As explained in the Introduction, we have also evaluated  
 116 the dose-average LET as a dosimetry parameter that may provide a more complete physical

TABLE I: Minimum and maximum energies used to generate SOBP for each ion.

	$^{12}\text{C}$	$^{16}\text{O}$	$^{20}\text{Ne}$	$^{28}\text{Si}$	$^{40}\text{Ar}$	$^{56}\text{Fe}$
Emin (MeVu $^{-1}$ )	164.6	195.5	221.2	270.4	297.5	382.8
E <sub>max</sub> (MeVu $^{-1}$ )	194.2	230.6	276.0	319.0	351.0	450.0

117 description of heavy ion SFRT. LET describes the energy transferred into a narrow region  
 118 around the primary ion track. While the LET of a pure beam of ions with a fixed energy  
 119 is well defined, the LET of a mixed radiation field is more complex. The reason for that is,  
 120 that in a mixed radiation field, LET has to be averaged over the different ions contributing.  
 121 This is often done by using the so-called dose-averaged LET, where the LET of each particle  
 122 is weighted according to the dose it is contributing. GATE v8.0 was used to calculate 3D  
 123 maps of unrestricted dose-averaged LET. In particular, we have evaluated the dose-averaged  
 124 LET in the peak and valley regions as a function of depth. We have also introduced a new  
 125 parameter, called Peak-to-valley-LET-ratio (PVL<sub>R</sub>), to give a more thorough dosimetric  
 126 quantification of this approach. Two different (*ctc*) distances have been considered: a *ctc*  
 127 leading to LET-homogeneous profiles in the target and a *ctc* that has been shown to minimize  
 128 the contribution of heavy nuclear fragments to the valleys, i.e. 3500  $\mu\text{m}$  [3] and [27].

### 129 III. RESULTS AND DISCUSSION

130 In this section, the results of the assessment of the minimum *bw* for each ion and config-  
 131 uration are reported first. Then, PVDR values as a function of depth corresponding to that  
 132 *bw* will be shown. Finally, dose-averaged LET distributions will be presented.

133 In order to assess the minimum *bw* for each ion for the two irradiation geometries under

TABLE II: Optimum beam width determined for each ion.

	$^{12}\text{C}$	$^{16}\text{O}$	$^{20}\text{Ne}$	$^{28}\text{Si}$	$^{40}\text{Ar}$	$^{56}\text{Fe}$
MBRT $bw$ ( $\mu\text{m}$ )	600	600	500	500	400	400
GRT $bw$ ( $\mu\text{m}$ )	900	900	800	800	700	700

134 consideration, MBRT and GRT, depth doses curves were evaluated following the dosimetry  
 135 criterium explained in section II. See figure2. The shape of the depth dose curves significantly  
 136 change when very small  $bw$  are used: a (very) high ratio of lateral scattered dose with respect  
 137 to the dose deposited in the primary trajectory of the beam, resulting in the loss of the  
 138 characteristic Bragg peak shape and its advantages. The same reason explains the difference  
 139 between MBRT and GRT configurations. Whether high entrance-to-Bragg peak dose ratio  
 140 and the relative higher presence of nuclear fragments in the valleys could be compensated in  
 141 terms of tissue sparing by distinct biological mechanisms being activated when very narrow  
 142 beams ( $\leq 100 \mu\text{m}$ ) are used is yet to be established by biological experiments. Table III  
 143 shows the results on the optimum (minimum)  $bw$  for each beam and configuration.

144 As expected the  $bw$  is smaller for the heaviest ions due to the reduced lateral scattering.  
 145 By using those minimum  $bw$  we have then assessed the PVDR, as well as LET maps. SOBP  
 146 have been used for that evaluation. Figure 3 shows examples of SOBP for Si ions for MBRT  
 147 and GRT configurations. The normal tissue at the entrance receive between 70-75 % in  
 148 MBRT and 65-70 % in GRT of the average dose in the SOBP for the different ions. Similar  
 149 values (70 %) have been reported in ‘standard’ ion therapy [37].

150 As explained in Materials and Methods, two different  $ctc$  were considered: a  $ctc$  that has  
 151 been shown to minimise the contribution of heavy nuclear fragments to the valleys, i.e., 3500

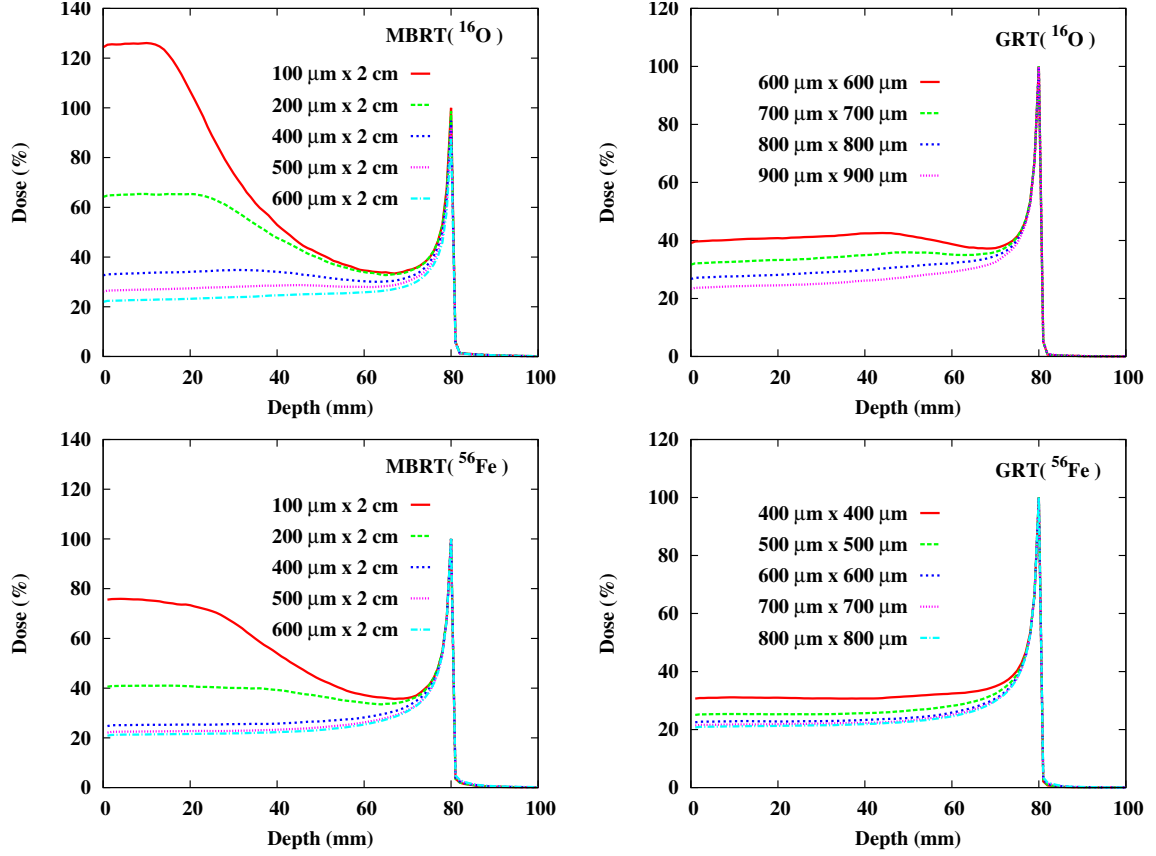


FIG. 2: Depth dose curves resulting from the irradiation (Pristine peaks) with beams of different widths. The upper row corresponds to curves of oxygen ions, the lower one to iron. The left and right columns depict the curves for MBRT and GRT irradiations, respectively. The values were normalized to the maximum value.

152  $\mu\text{m}$  [3] and [27] and a *ctc* leading to LET-homogenous profiles in the target. Concerning the  
 153 latter, an empirical relation between the *ctc* and the *bw* was found:

$$ctc = 2 \times bw, \quad \text{for MBRT} \quad (1)$$

$$ctc = 2 \times bw - 400 \mu\text{m}, \quad \text{for GRT} \quad (2)$$

154 Figure 4 shows the PVDR for the different ions, configurations and *ctc* distances. Ex-  
 155 tremely high PVDR values for a *ctc* of 3500  $\mu\text{m}$  are obtained, similar or equal to previously

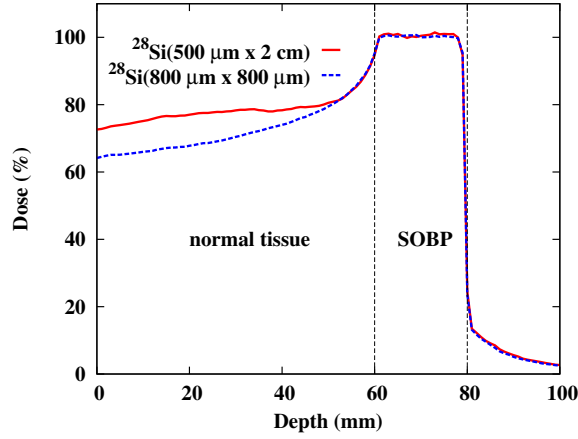


FIG. 3: Percentage depth dose curves for the central peak silicon MBRT (solid line) and GRT (dotted line). The values were normalized with respect to the average dose in the SOBP.

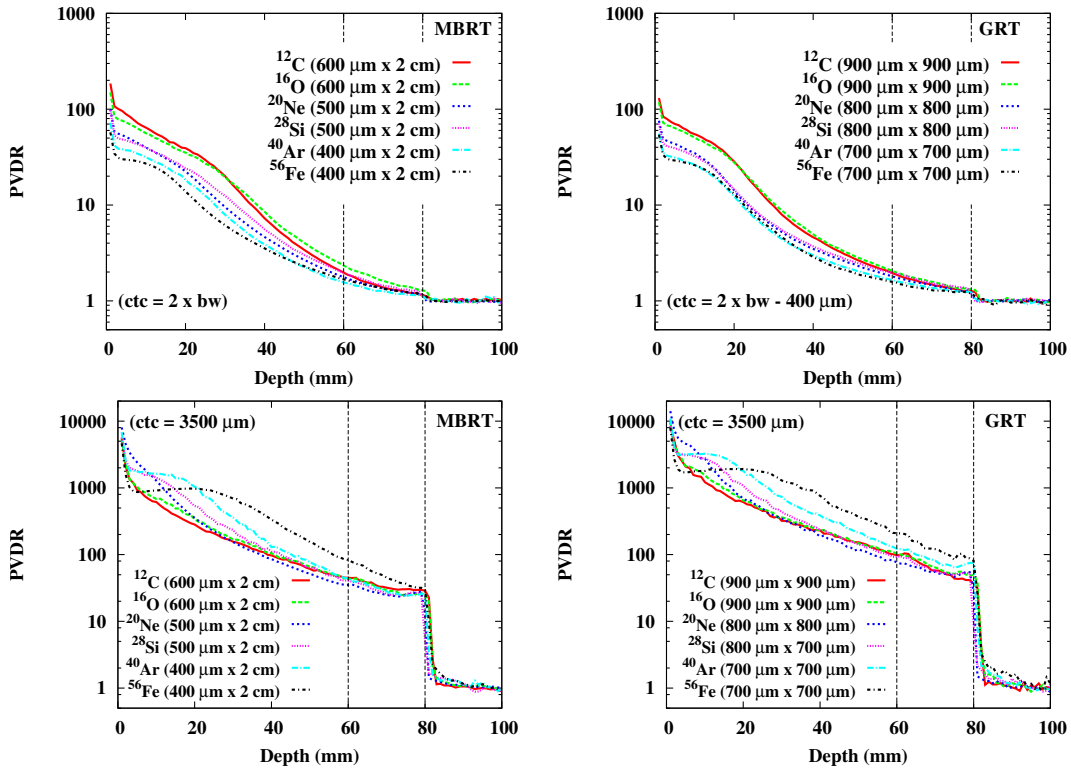


FIG. 4: PVDR as a function of depth for each ion. The left and right columns corresponds to MBRT and GRT, respectively. The upper row shows the results when the *ctc* leading to in the LET-homogenous profiles in the target is used, the lower row depicts the values for a *ctc* of 3500  $\mu\text{m}$ .

156 published works [3] and [27]. For the narrowest *ctc* distances, PVDR values in the first 2

157 cm are similar to those obtained in x-rays and proton MBRT [22], [38]. Values comparable  
158 to GRID therapy are achieved for depths from 2 to 5 cm [39]. A quasi homogeneous dose  
159 distribution is obtained in the target (from 6 to 8 cm depth) with only one array, without  
160 interlacing.

161 The narrower *ctc* leads to similar values for each ion in MBRT and GRT configurations.  
162 The lowest PVDR are obtained for the heaviest ions. The opposite situation is found when  
163 a *ctc* of 3500  $\mu\text{m}$  is used. Larger PVDR are obtained in GRT than in MBRT when a *ctc* of  
164 3500  $\mu\text{m}$  is used.

165 Figure 5 shows the ratio between the peak and valley dose-averaged LET (PVL<sub>R</sub>). The  
166 highest values are obtained in the first millimeters, being Fe the ion providing the highest  
167 PVL<sub>R</sub> ( $205 \pm 2$  in MBRT and  $200 \pm 2$  in GRT). In normal tissues in the beam path, the  
168 heaviest the ion, the highest the PVL<sub>R</sub>. The ratio between the PVL<sub>R</sub> in the entrance and  
169 the target region being as well the highest for the heaviest ion. A factor 100 is obtained  
170 in the entrance between the PVL<sub>R</sub> for Fe and C. PVL<sub>R</sub> as a function of depth for lightest  
171 ions, like carbon and oxygen have a soft trend, being in the entrance approximately double  
172 the value in the target. The curves of PVL<sub>R</sub> as a function of depth are very similar in the  
173 two configurations, MBRT and GRT, being slightly higher for GRT. As opposed to PVDR,  
174 PVL<sub>R</sub> does not significantly change with *ctc*. PVL<sub>R</sub> in the target is not homogeneous to  
175 the 3500  $\mu\text{m}$  *ctc* setup. The trend of the curves can be explaining using the dose-averaged  
176 LET in the valleys (figure 7), to this setup, the trend of the curves are different that the  
177 trend of the curves to the small *ctc* setup.

178 Figures 6 and 7 show the dose-averaged LET in the peaks and the valleys, respectively.  
179 Concerning the LET in the peaks, there are not significant differences between MBRT and  
180 GRT, not variation with the *ctc*. In contrast, LET increases with the atomic number of the

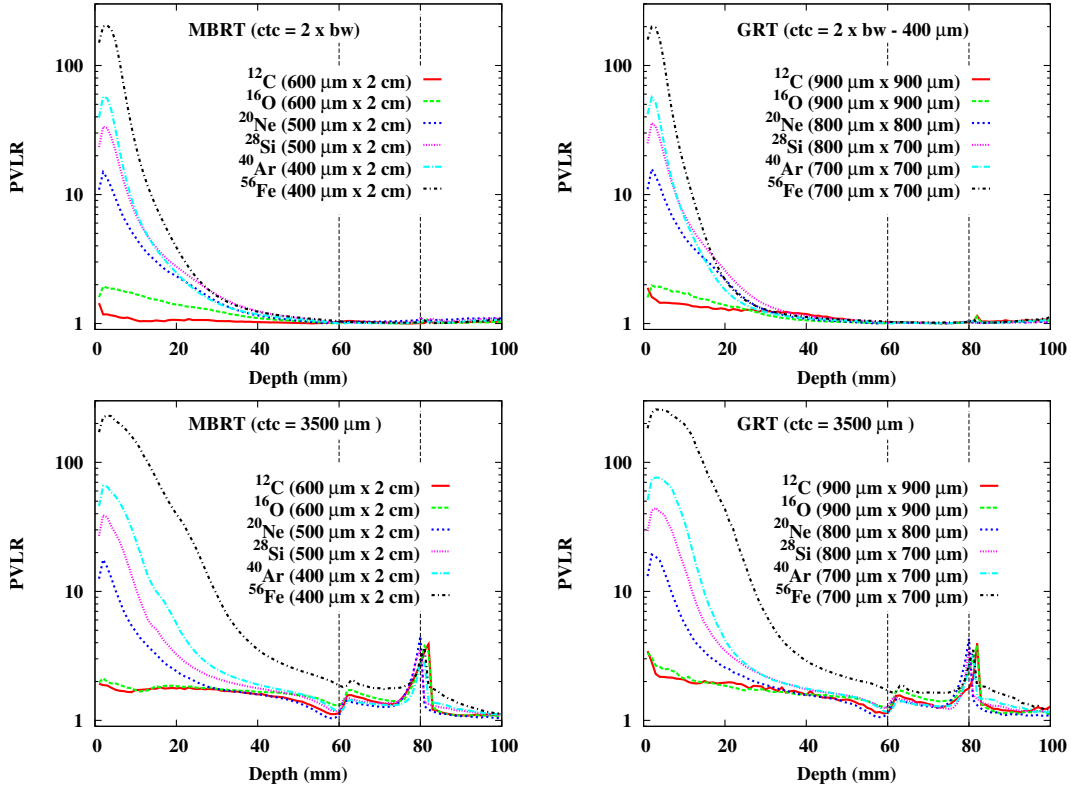


FIG. 5: PVLr as a function of depth for each ion. The left and right columns corresponds to MBRT and GRT, respectively. The upper row shows the results when the ctc leading to in the LET-homogenous profiles in the target is used, the lower row depicts the values for a *ctc* of 3500  $\mu\text{m}$ .

181 ion, as expected.

182 In ‘standard’ hadron therapy (non spatially fractionated), LET values of about 200  
 183  $\text{keV}/\mu\text{m}$  produce the largest biological effectiveness [40]. An “overkill” effect is reported  
 184 for LET higher than 200  $\text{keV}/\mu\text{m}$  for  $^{12}\text{C}$  ions. Some works [41] indicate that the exclu-  
 185 sion of non-hit fraction in the calculation of surviving fraction partially prevented the fall  
 186 of biological effectiveness when LET exceeded 200  $\text{keV}/\mu\text{m}$ . The presence of non-hit cells  
 187 among hit cells is one of the features characteristic for exposure to very high-LET heavy  
 188 ions. Mehnati *et al.* also suggest that bystander effect may occur in non-hit cell. The  
 189 expected radiobiological features of minibeam radiation therapy are so different from those

190 of ‘standard’ RT, including non-targeted effects (see Y. Prezado *et al.* 2015), that the ra-  
191 diobiological effect cannot be assessed by means of established RBEs of ‘standard’ hadron  
192 therapy (broad beam irradiation). It could be that the fact of having more severe damage  
193 in fewer sites triggers a more important contribution of non-targeted effects. Whether in  
194 these conditions a possible overkill effect in the distal part of the Bragg peak would be  
195 totally maintained or compensated by other phenomena appearing in spatially fractionated  
196 techniques, it is yet to be explored.

197 Values lower than  $200 \text{ keV}/\mu\text{m}$  are found for all ions except Fe and the distal part of the  
198 SOBPs of Si and Ar. Whether an overkill effect of the same magnitude than in ‘standard  
199 RT’ would start participating in this technique is yet to be assessed.

200 Valleys are believed to be (main) responsible for tissue sparing [18]. Low LET values in  
201 the valleys could then favor normal tissue preservation. The shape of the LET curves is  
202 similar for MBRT and GRT. A smooth variation of the dose-averaged LET as a function of  
203 depth is obtained for carbon and oxygen ions. The values for carbon are lower than those  
204 encountered in clinics. In contrast, a continuous increase is observed for the heavier ions,  
205 with a very high gradient in the first centimeters. An interesting feature is the fact that the  
206 LET values for Ne and heavier ions are lower than those of carbon ions in the first centimeter,  
207 which would favor normal tissue preservation. The depth where the trend is reversed depends  
208 on *ctc* (the larger the *ctc*, the deeper the point) and on ion (the heavier the ion, the deeper  
209 the point). The reason is the heavier the primary ion is, the higher the proportion of  
210 heavier fragments (high LET), whose angular distribution, mainly determined by reaction  
211 kinematics, will then be more forward directed than lighter products. Consequently, their  
212 dose deposition will occur deeper. In normal tissues, all ions but Fe, lead to LET values  
213 lower than  $100 \text{ KeV}/\mu\text{m}$ . Superior values are considered to lead to an RBE higher than one

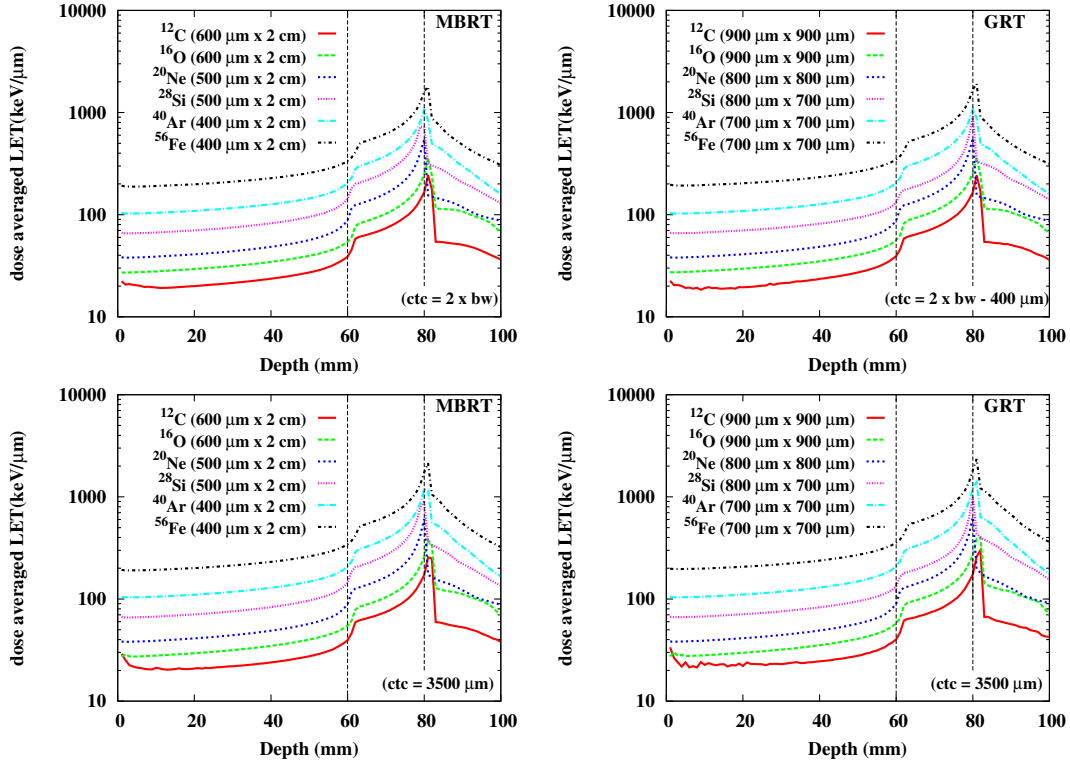


FIG. 6: Dose-averaged LET (peak region) as a function of depth for the different ions. The left and right column corresponds to MBRT and GRT irradiations, respectively. The upper row shows the results when the *ctc* leading to in the LET-homogenous profiles in the target is used, the lower row depicts the values for a *ctc* of 3500  $\mu\text{m}$ .

214 [42]. In the target area, Ne and lightest ions provide LET values lower than 200 keV/ $\mu\text{m}$  at  
 215 all depths. LET values are a factor 2 smaller when the 3500  $\mu\text{m}$  *ctc* is used. Not significant  
 216 differences are observed between MBRT and GRT configurations.

217 Among the different ions, Ne could offer a good compromise since they lead to high  
 218 PVDR and PVLr in normal tissues, LET values in the valleys lower than carbon in the first  
 219 centimeter, remaining below 30 keV/ $\mu\text{m}$  in normal tissues (up 6 cm depth) and providing  
 220 values close to 100 keV/ $\mu\text{m}$  in the target region. In addition, Ne provides a lower oxygen  
 221 enhancement ratio (OER) than lighter ions, which is advantageous for the treatment of

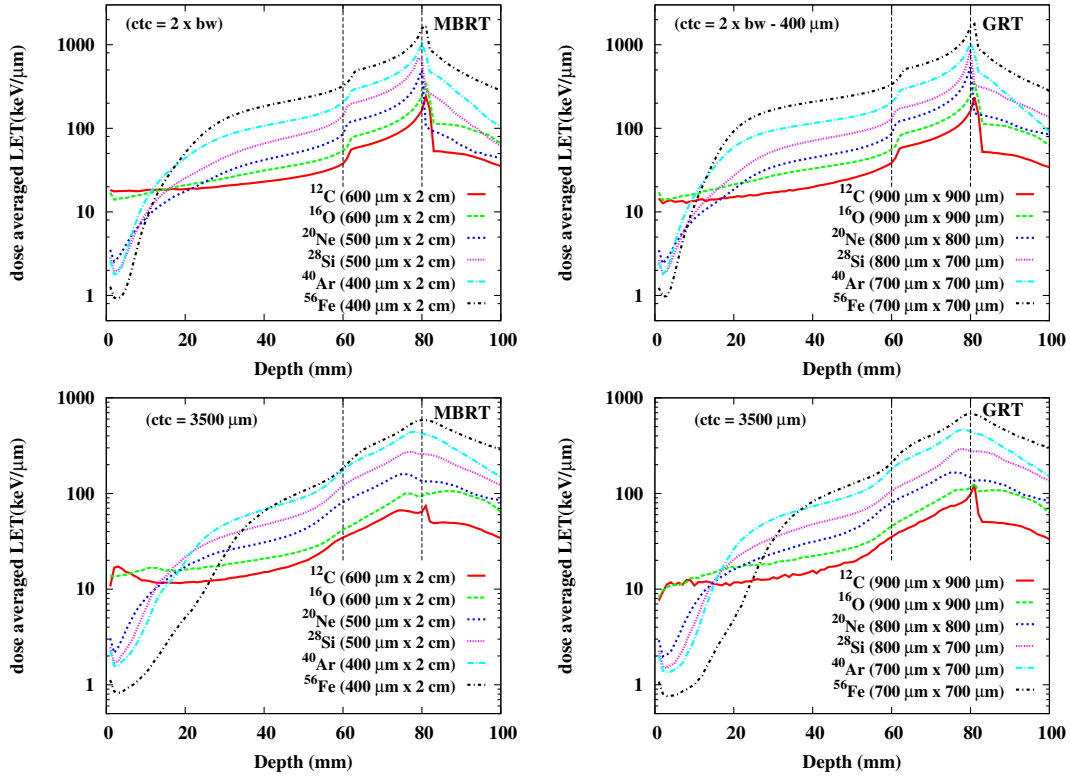


FIG. 7: Dose-averaged LET (valley region) as a function of depth for the different ions. The left and right column corresponds to MBRT and GRT irradiations, respectively. The upper row shows the results when the *ctc* leading to in the LET-homogenous profiles in the target is used, the lower row depicts the values for a *ctc* of 3500  $\mu\text{m}$ .

222 hypoxic tumors.

223 The larger *ctc* leads to higher PVDR and reduced valley doses, with a lower contribution  
 224 of heavier nuclear fragments to the valleys [3], [27] leading to a factor 2 decrease in the dose  
 225 averaged LET in normal tissues. The PVL<sub>R</sub> in the target is lower than 2 (except for Fe),  
 226 so quasi homogeneous.

227 Concerning the differences between GRT and MBRT, the main difference is that GRT  
 228 leads to higher PVDR than MBRT a factor 1.6, however the minimum beam width 1.5 times  
 229 larger, which might reduce the tissue tolerances.

230 Although only biological experiments can establish the most favorable configuration,  
 231 the results obtained in the physical study suggest that Ne MBRT with beam width of  
 232  $500 \mu\text{m} \times 2 \text{ cm}$  and a *ctc* of  $3500 \mu\text{m}$  could be the best candidate. Figure 8 shows a 2D dose  
 233 map for this case.

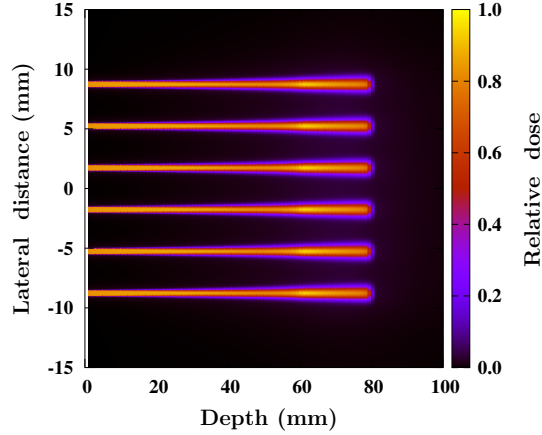


FIG. 8: 2D dose map for the MBRT with beam width of  $500 \mu\text{m} \times 2 \text{ cm}$  and a *ctc* of  $3500 \mu\text{m}$  .

#### 234 IV. CONCLUSIONS

235 Heavy ions minibeam or Grid therapy is a new approach for the treatment of radiore-  
 236 sistant tumors. The normal tissue preservation that could be offer by this strategy would  
 237 allow a dose escalation in the tumor, and therefore, a higher tumor control may be expected.  
 238 This work aimed at the optimization of the irradiation configuration in terms of ion species,  
 239 beam width, beam spacing and LET-maps. The use of submillimetric beams still allows to  
 240 keep a ratio between the dose in the entrance and the dose in the target of the same order as  
 241 in conventional irradiations. MBRT permits the employment of thinner beams than GRT.  
 242 Large *ctc* distances ( $3500 \mu\text{m}$ ) should be preferred since they lead to valley doses composed  
 243 of lighter nuclear fragments resulting in a much reduced dose-averaged LET values in normal  
 244 tissue, favoring its preservation. However, preclinical studies evaluating tumor response will

245 be the ones providing the ultimate determination of the most effective etc distance for tumor  
246 control. Among the different ions species evaluated, Ne stands out as the one leading to the  
247 best balance between high PVDR and PVLR in normal tissues and high LET values (close  
248 to  $100 \text{ keV}/\mu\text{m}$ ) and a favorable OER in the target region. Although biological experiments  
249 are warranted to confirm these conclusions, the results of this work may be used to guide  
250 the future biological studies.

### 251 **Acknowledgment**

252 The authors acknowledge the calculation time granted at Centre de Calcul de Lyon  
253 (IN2P3) and at Grand Equipement National de Calcul Intensif (in particular at the super-  
254 computer Curie of CEA). We warmly thank the access to Castilla y León Supercomputing  
255 Center (Calendula) awarded by the Red Española de Supercomputación. We also acknowl-  
256 edge PRACE for awarding us access to resource MareNostrum Barcelona Supercomputing  
257 Center based in Spain.

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