

Energy-Transfer Sensitization of BODIPY Fluorescence in a Dyad with a Two-Photon Absorbing Antenna Chromophore

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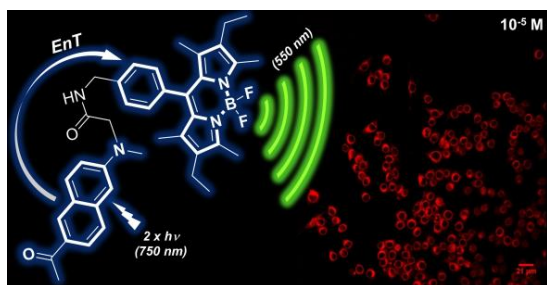
Highlights

- A BODIPY-Prodan dyad with a light-harvesting antenna and a non-conjugating linker was synthesized.
- Highly efficient two-photon-induced energy transfer mediates the excited state communication between the two chromophores.
- Intramolecular photoinduced electron transfer is controlled by solvent polarity.
- Two-photon excitation under conditions of the imaging of N13 microglial cells is demonstrated.

Abstract

A bichromophoric dyad composed of a Prodan-derived moiety and a BODIPY dye was prepared, photophysically characterized, and tested for multiphoton fluorescence microscopy of N13 microglial cells. The dyad shows highly efficient energy transfer, independent on the solvent polarity of the medium. However, in polar media photoinduced electron transfer is a competitive pathway. The Prodan-derived chromophore shows significant two-photon absorption. Hence, it can be used a two-photon-absorbing antenna for the sensitization of BODIPY fluorescence.

Graphical abstract



Keywords

BODIPY dyes

Energy transfer

Electron transfer

Two-photon absorption

Bioimaging

1. Introduction

Boron dipyrromethene (BODIPY) compounds are among the archetypal functional dyes in bio-relevant applications and have found widespread use as labels for cell imaging [1–7], temperature or viscosity probes [8–10], chemosensors [11–14], and photosensitizers for singlet oxygen [4,15–19]. The preference as an imaging probe is motivated by the high molar absorption coefficients (*ca.* 60000–80000 M⁻¹cm⁻¹), paired with highly efficient fluorescence (emission quantum yields Φ_f *ca.* 0.6–1.0) [17,20–24]. A generally observed drawback of BODIPY dyes is their rather small Stokes shift (*ca.* 500–800 cm⁻¹), which may cause problems such as excitation straylight detection, re-absorption of emitted photons or homo-energy transfer [23–25]. A viable strategy to overcome this shortcoming is the structural integration in antenna-substituted dyad designs that enable the observation of the highly efficient BODIPY fluorescence emission with a virtually increased Stokes shift by means of through-space or through-bond energy transfer (EnT) [26–30]. On the one hand, in the latter case electronic communication between the energy-donor antenna and the BODIPY energy acceptor is facilitated in so-called EnT cassettes [12,26–28,31–33]. On the other hand, through-space EnT is classically described by the Förster theory, implying antenna-BODIPY dyads with non-conjugating linkers [29,34–37].

A second strategy which eliminates the mentioned difficulties is the design of two-photon-excitable BODIPY chromophores [3,38–40]. For this purpose, BODIPY dyes are modified with electron-donating aromatic amines and conjugating ethenyl linkers, thereby enabling electronic communication [3,40–42]. These probes were successfully exploited for two-photon-excitation (TPE) fluorescence microscopy for cellular imaging [3,43]. The use of red or near-infrared excitation light sources bears additional

advantages such as deeper penetration into biological tissues, less background fluorescence from biomolecules, and improved photostability of the dyes [23,44,45]. In the light of these two approaches for improved BODIPY dyes in the context of biologically relevant applications we sought to synthesize and characterize a dyad that combines both strategies by means of using a two-photon-excitable antenna to sensitize efficiently BODIPY fluorescence by EnT. Two-photon-induced EnT is less frequently reported and requires a careful selection of the antenna energy donor dye [29,46–48]. Herein we selected a fluorophore class with significant intramolecular charge-transfer (ICT) properties, known to exhibit significant two-photon-absorption cross sections (σ_2): a 2-*N,N*-dialkylamino-6-alkanoylnaphthalene derivative. The latter motif is found in archetypal environment-sensitive fluoroprobes such as Prodan (2-dimethylamino-6-propionynaphthalene) [49], Badan (6-bromoacetyl-2-dimethylaminonaphthalene) [50] or Laurdan (2-dimethylamino-6-lauroynaphthalene) [51]. These dyes have enjoyed extensive application for the monitoring of the biophysical properties of lipid membranes, including parameters such as polarity, viscosity, and structural order [52–54].

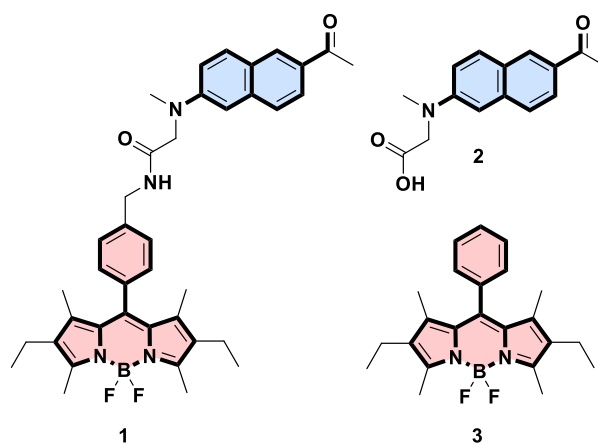


Fig. 1. Structures of the bichromophoric dyad **1** and the individual model chromophores **2** and **3**.

The concrete structures of the herein prepared bichromophoric antenna-BODIPY dyad and the separate chromophore models are shown in Fig. 1. We present their complete photophysical characterization, including the highly efficient EnT-sensitization of the BODIPY fluorescence, the implications of solvatochromic effects on the antenna excited-state properties, and the two-photon-absorption processes. Finally, we demonstrate the internalization of the dyad **1** in N13 microglial cells and proof-of-concept for the validity of the molecular and photophysical design in complex bio-relevant environments.

2. Materials and methods

2.1. General

All reagents and solvents were obtained from commercial sources and were used without further purification unless otherwise stated. Tetrahydrofuran (THF), acetonitrile (MeCN), *N,N*-dimethylformamide (DMF), and methanol (MeOH), which were used in the photophysical measurements, were of spectroscopy grade. ¹H, ¹³C, ¹¹B, and ¹⁹F NMR spectra were recorded either on a Bruker Advance III 400 MHz or a Bruker Advance III 500 MHz instrument. Structural assignments were made with additional information from gCOSY and gHSQC experiments. Chloroform-*d* and dimethylsulfoxide-*d*₆ (all 99 atom-% D) were used as solvents, and the spectra were referenced to the residual solvent peak ($\delta = 7.26$ and 2.50 ppm, respectively). High-resolution mass spectrometry (HRMS) was performed using an Elite QTOF, Bruker Daltonics Autoflex ESI-TOF or a Thermo Fisher Orbitrap high-resolution mass spectrometer. Fourier transform infrared spectra (FTIR) were collected on a Bruker Tensor27 spectrometer or an FT/IR 4200 Jasco spectrometer with an attenuated total reflectance (ATR) module. The samples were supported in KBr and the spectra were

recorded in the 500–4000 cm^{-1} wavenumber range. The products were purified by column chromatography or by preparative thin-layer chromatography on silica gel Merck-60 (230-400 mesh, 60 Å) and analytical TLC was performed on aluminum sheets pre-coated with silica gel 60 F₂₅₄ (Merck).

2.2. *Synthesis and characterization*

4-((2,2,2-Trifluoroacetamido)methyl)benzoic acid **4** [55] (Scheme 1) and the model chromophores Prodan **2** [52] and BODIPY **3** [56] were prepared following reported procedures. The analytical characterization of the final products **2** and **3** revealed their identity and purity (Supplementary Material – Fig. S7–S14, S22–S23, and S27–S28).

2.2.1. *BODIPY 5*

An oven-dried two-neck round-bottom flask, connected to a reflux condenser and held under a nitrogen atmosphere, was charged with the protected amino derivative **4** (150.0 mg, 0.61 mmol). Toluene (10.0 mL) and thionyl chloride (1.0 mL, 13.7 mmol) were added, and the reaction mixture was stirred at 80 °C for 12 h. The solvent was evaporated under reduced pressure and the residue was dissolved in dry dichloromethane (20 mL). Then kryptopyrrole (0.3 mL, 2.2 mmol) was added and the mixture was refluxed for 4 hours under a nitrogen atmosphere. Subsequently, triethylamine (2.0 mL) and $\text{BF}_3 \cdot \text{OEt}_2$ (2.0 mL) were added into the mixture upon which it turned dark red. The reaction mixture was then diluted with dichloromethane (100 mL) and washed with water (2 × 20 mL). The organic layer was dried over anhydrous MgSO_4 and the solvent was evaporated under reduced pressure. The crude was purified by silica gel column chromatography (dichloromethane/cyclohexane, 1:1 to pure dichloromethane gradient, $R_F = 0.56$) affording BODIPY **5** (61 mg, 20% yield) as a red

amorphous solid. ^1H NMR (400 MHz, CDCl_3 , 298 K): δ 7.40 (d, $J = 8.0$ Hz, 2H), 7.29 (d, $J = 8.0$ Hz, 2H), 6.90 (br, 1H), 4.63 (d, $J = 5.9$ Hz, 2H), 2.52 (s, 6H), 2.30 (q, $J = 7.5$ Hz, 4H), 1.26 (s, 6H), 0.97 (t, $J = 7.5$ Hz, 6H) ppm (Fig. S15). $^{13}\text{C}\{^1\text{H}\}$ NMR (100 MHz, CDCl_3 , 298 K): δ 158.0, 157.6, 157.2, 156.8, 154.0, 139.3, 138.2, 136.9, 135.8, 133.0, 130.7, 129.0, 128.2, 120.2, 117.3, 114.5, 111.6, 43.4, 17.1, 14.6, 12.5, 11.8 ppm (Fig. S16). HRMS (ESI): Found $[\text{M}+\text{H}]^+$ 506.2408; molecular formula $\text{C}_{26}\text{H}_{30}\text{BF}_5\text{N}_3\text{O}$ requires $[\text{M}+\text{H}]^+$ 506.2397 (Fig. S24). FTIR (ATR) ν_{max} : 3390, 3373, 2965, 2930, 2872, 1727, 1538, 1321, 1190, 1161, 1063, 974, 798, 713, 537 cm^{-1} (Fig. S29).

2.2.2. BODIPY 6

An oven-dried 25 mL Schlenk tube was charged with BODIPY 5 (60.0 mg, 0.12 mmol), anhydrous K_2CO_3 (20.0 mg, 0.14 mmol), and 10 mL of a methanol/water mixture (9:1, 10 mL). The solution was stirred at room temperature for 48 hours. The mixture was then concentrated under reduced pressure, re-dissolved in dichloromethane (20 mL), and washed with water (2×20 mL). The organic layer was dried over anhydrous MgSO_4 and the solvent was evaporated under reduced pressure to afford a red amorphous solid (38 mg, 78% yield) which was used in the next step without further purification. ^1H NMR (400 MHz, CDCl_3 , 298 K): δ 7.44 (d, $J = 7.5$ Hz, 2H), 7.25 (d, $J = 7.5$ Hz, 2H), 3.99 (s, 2H), 2.54 (s, 6H), 2.31 (q, $J = 7.2$ Hz, 4H), 1.30 (s, 6H), 0.99 (t, $J = 7.2$ Hz, 6H) ppm (Fig. S17). $^{13}\text{C}\{^1\text{H}\}$ NMR (100 MHz, CDCl_3 , 298 K): δ 153.5, 143.8, 140.2, 138.3, 134.1, 132.6, 130.8, 128.3, 127.5, 46.0, 29.6, 17.0, 14.5, 12.4, 11.7 ppm (Fig. S18). HRMS (ESI): Found $[\text{M}+\text{H}]^+$ 410.2583; molecular formula $\text{C}_{24}\text{H}_{31}\text{BF}_2\text{N}_3$ requires $[\text{M}+\text{H}]^+$ 410.2574 (Fig. S25). FTIR (ATR) ν_{max} : 2964, 2928, 2870, 1671, 1538, 1474, 1319, 1265, 1187, 1061, 975, 758, 710 cm^{-1} (Fig. S30).

2.2.3. Dyad 1

An oven-dried 25 mL Schlenk tube was charged with the Prodan derivative **2** (30.0 mg, 0.12 mmol), benzotriazole-1-yl-oxy-tris-pyrrolidino-phosphonium hexafluorophosphate (PyBOP; 62.0 mg, 0.12 mmol), 4-methylmorpholine (NMM; 0.15 mL, 1.37 mmol), and dry dichloromethane (5.0 mL) under nitrogen atmosphere. The mixture was stirred at room temperature for 30 minutes. Subsequently, BODIPY **6** (48 mg, 0.12 mmol) was added in several portions and the reaction mixture was stirred for another 90 minutes at room temperature. The mixture was then diluted with dichloromethane (100 mL). The organic phase was washed with milli-Q water (2×20 mL), followed by 0.1 M HCl (2×20 mL). The organic layer was dried over anhydrous MgSO_4 and the solvent was evaporated under reduced pressure. The crude product was purified by silica gel column chromatography (ethyl acetate/*n*-hexane 4:6), followed by preparative TLC (dichloromethane/methanol 50:1, $R_F = 0.69$), affording the desired dyad **1** (31 mg, 41%) as red amorphous solid. ^1H NMR (500 MHz, CDCl_3 , 298 K): δ 8.33 (s, 1H), 7.96 (dd, $J = 8.7$, 1H), 7.85 (d, $J = 9.1$ Hz, 1H), 7.68 (d, $J = 8.7$ Hz, 1H), 7.29 (d, $J = 8.0$ Hz, 2H), 7.15 (d, $J = 8.0$ Hz, 2H), 7.11 (dd, $J = 9.1$, 2.6 Hz, 1H), 7.00 (s, 1H), 6.83 (t, $J = 6.1$ Hz, 1H), 4.60 (d, $J = 6.1$ Hz, 2H), 4.14 (s, 2H), 3.21 (s, 3H), 2.66 (s, 3H), 2.51 (s, 6H), 2.25 (q, $J = 7.5$ Hz, 4H), 1.11 (s, 6H), 0.95 (t, $J = 7.5$ Hz, 6H) ppm (Fig. S1). $^{13}\text{C}\{^1\text{H}\}$ NMR (126 MHz, CDCl_3 , 298 K): δ 197.7, 169.9, 154.0, 148.8, 139.7, 139.0, 138.3, 137.3, 135.2, 132.9, 132.2, 131.4, 130.8, 130.1, 128.8, 128.0, 126.7, 126.4, 125.3, 116.2, 107.1, 58.5, 42.9, 40.2, 26.6, 17.2, 14.7, 12.6, 11.8 ppm (Fig. S3). ^{11}B NMR (160 MHz, CDCl_3 , 298 K): δ 0.77 (t, $J = 33.6$ Hz) ppm (Fig. S5). ^{19}F NMR (470 MHz, CDCl_3 , 298 K): δ -145.81 (q, $J = 32.4$ Hz) ppm (Fig. S6). HRMS (ESI): Found $[\text{M}+\text{Na}]^+$ 671.3395; molecular formula $\text{C}_{39}\text{H}_{43}\text{BF}_2\text{N}_4\text{O}_2$ requires $[\text{M}+\text{Na}]^+$ 671.3346 (Fig. S20). FTIR (KBr)

ν_{max} : 3413, 2959, 2924, 2855, 1665, 1621, 1461, 1378, 1261, 1095, 1021, 801 cm^{-1}

(Fig. S26).

2.3. Photophysical measurements

All photophysical measurements were performed at room temperature (293 K) with 10 μM solutions, contained in quartz cuvettes (1 cm optical pathlength). UV/Vis absorption spectra were measured with a UVPC-1603 instrument from Shimadzu. Steady-state fluorescence measurements were done on a Varian Eclipse fluorimeter. The fluorescence spectra were corrected for the spectral responsivity of the photomultiplier-based detection system by applying an emission correction curve. Fluorescence quantum yields were determined for diluted sample solutions (absorbance of ≤ 0.1 at the excitation wavelength) and by using rhodamine 6G in ethanol ($\Phi_{\text{f}} = 0.95$) [57] or quinine sulfate in 0.05 M H_2SO_4 as reference ($\Phi_{\text{f}} = 0.55$) [58]. The quantum yields were calculated by taking the refractive indices of the employed solvents into account [59]. The two-photon absorption cross-sections (σ) were determined according to a standard two-photon excited-fluorescence method [60]. The experiments were done within a laser-power regime where the fluorescence is proportional to the square of the laser excitation power. As reference a rhodamine B solution (10 μM in MeOH) was used [60], reasonably assuming that the fluorescence quantum yield is the same regardless of one- or two- photon excitation. The two-photon-excited fluorescence of dyad **1** was analysed using an inverted Leica SP5 MP confocal microscope equipped with a MaiTai Ti:Sapphire HP laser (Spectra-Physics, Inc.), tuneable between 690 and 1040 nm. Emission spectra were measured using a dynamic 10 nm wide emission detection window moving in 28 steps between 425 and 705 nm.

Time-correlated single-photon-counting measurements were performed on a FS920 fluorescence spectrometer (Edinburgh Instruments). As excitation sources, picosecond pulsed diode lasers (Edinburgh Instruments) were used (EPL-445: $\lambda_{\text{exc}} = 442.2$ nm, pulse width 78.3 ps; EPL-485: $\lambda_{\text{exc}} = 482.0$ nm, pulse width 100.6 ps). The nanosecond fluorescence decays were deconvoluted, using the instrument response function that was recorded with a light-scattering Ludox solution.

2.4. DFT calculations

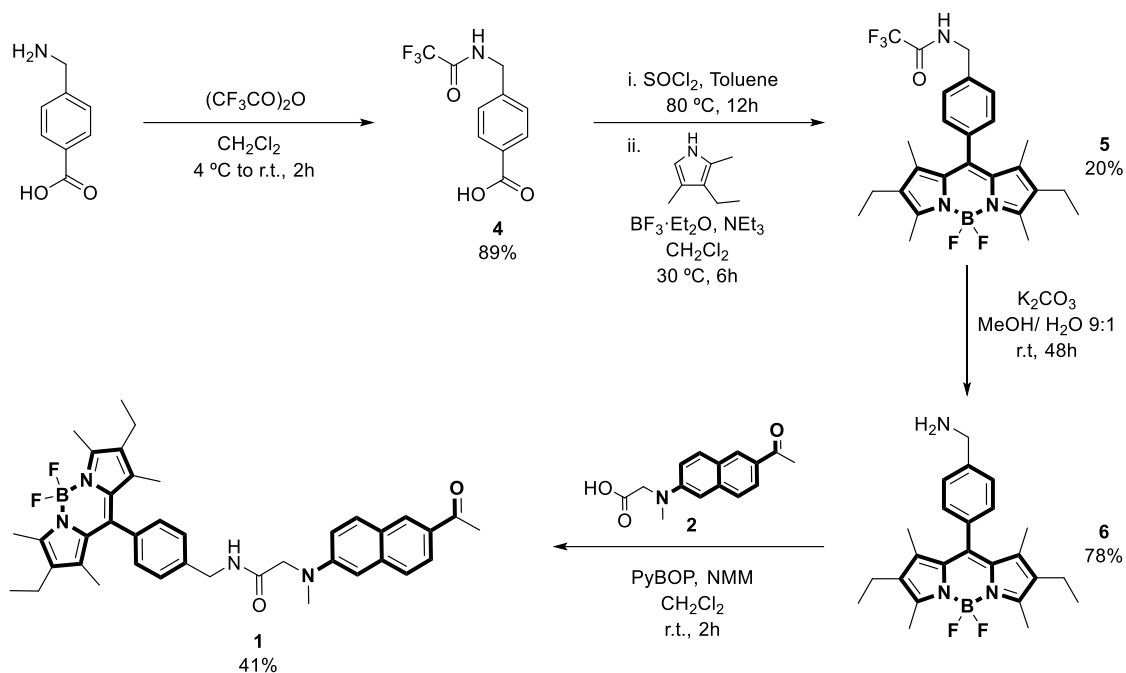
All calculations were done with the Gaussian 09 package [61]. The solvent effect was considered by including the polarizable continuum model (PCM) [62] and MeCN as the solvent. The starting geometries for the calculations were first optimized using AM1 followed by density functional theory (DFT). The ground-state (S_0) geometrical parameters have been determined with the density functional level of theory, without symmetry constraints and employing the Coulomb-attenuated CAM-B3LYP functional and the 6-31+G(d) basis set. Using the optimized ground-state geometries as starting coordinates, time-dependent density-functional-theory (TD-DFT) geometry optimizations of the (S_1) states were performed, using the first ten excited states and the linear-response (LR-PCM) approach. On these optimized structures, the state-specific model (SS-PCM) has been evaluated considering an equilibrated state between the solvent and the excited-state electron density.

The orientation factor (κ^2) is defined by $\kappa^2 = (\cos\theta_T - 3\cos\theta_D\cos\theta_A)^2$, where θ_T is the angle between the emission transition dipole of the donor and the absorption transition dipole of the acceptor, θ_D and θ_A are the angles between these dipoles and the vector joining the donor and acceptor [59]. The calculation of the orientation factor requires the transition dipole moment vectors of the donor and acceptor moieties, which

were obtained by carrying out TD-DFT (CAM-B3LYP/6-31+G(d) level of theory, first ten excited states). Once the transition electric dipole moment vector coordinates were obtained, these coordinates were appended on the optimized geometry of the molecule. The centers of the individual donor and acceptor moieties were joined and all the angles that are necessary to calculate the orientation factor were obtained.

2.5. *Bioimaging experiments*

N13 microglial cells were grown to 50–70 % confluence on 8-well slides at 37 °C with 5% CO₂ in Gibco RPMI 1640 medium supplemented with *L*-glutamine and fetal bovine serum. The N13 microglial cells were incubated for 1 hour with 1 mL of fresh, preheated RPMI 1640 culture medium containing varying concentrations (100 μM, 10 μM, 1 μM) of dyad **1** prepared from a stock solution ([**1**] = 2.1 mM) in DMSO. The intracellular fluorescence of dyad **1** was analysed using an inverted Leica SP5 MP and a HCX IRAPO L 25.0 × 0.95 WATER objective. Optimal conditions were maintained throughout the imaging process with an integrated microscope enclosure keeping cells at 37 °C and 5% CO₂. Dyad **1** was visualized using two-photon excitation at 750 nm and fluorescence detection between 425–705 nm using a HyD non-descanned detector. Brightfield images were captured independently using a 405-nm laser as light source.



Scheme 1. Synthesis of dyad **1**.

3. Results and discussion

3.1. Synthesis of dyad **1**

The bichromophoric dyad **1** was synthesized in a three-step sequence (Scheme 1) starting by preparation of the trifluoroacetyl-protected amino-BODIPY derivative **5** by means of an activation of the previously described protected *p*-aminobenzoic acid derivative **4** [55] in form of its acyl chloride (reaction with SOCl_2), reaction with kryptopyrrole, and finally treatment with $\text{BF}_3 \cdot \text{OEt}_2$ and NEt_3 as base. In the next step, the deprotection of the amino group was carried out under basic conditions (K_2CO_3) to obtain the BODIPY derivative **6**. In the final step **6** was then reacted with the known Prodan-derived antenna chromophore **2** [52] by means of a standard amidation protocol with the coupling reagents 4-methylmorpholine and PyBOP. This yielded the bichromophoric dyad **1**, containing a hydrolytically stable amide linker. Column chromatography on silica gel, followed by a second repurification by preparative TLC, yielded an analytical pure sample (see the Supplementary Material for details on the

detailed 1D and 2D NMR-spectroscopic characterization, signal assignments, and high-resolution mass spectrometry data).

3.2. *Optical spectroscopy and excited-state processes on one-photon excitation (OPE)*

In order to account for the known solvatochromic effects of Prodan-like chromophores, such as the herein employed antenna, we used solvents of varying polarity in the present study: acetonitrile (MeCN), tetrahydrofuran (THF), *N,N*-dimethylformamide (DMF), and methanol (MeOH). All photophysical data are summarized in Table 1. The normalized UV/Vis-absorption and emission spectra of the dyad **1** and the model chromophores (**2** and **3**) in MeCN are shown in Fig. 2 and Fig. 3, respectively (see the Supplementary Material for the other solvents – Fig. S31–S32).

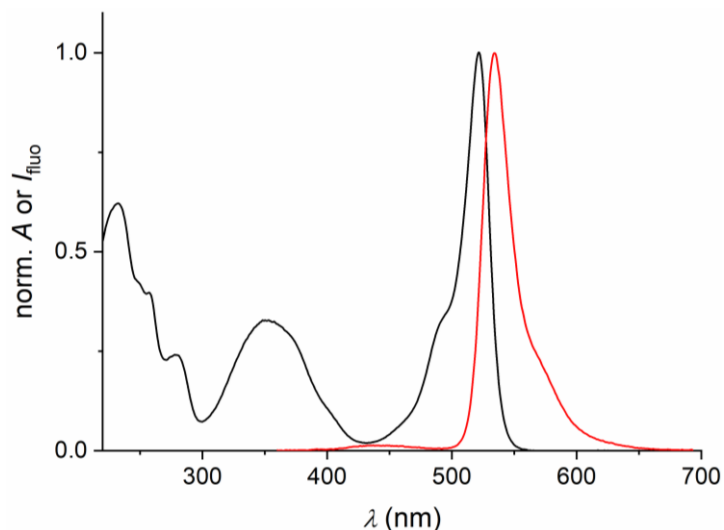


Fig. 2. UV/Vis absorption (black line) and fluorescence (red line; excitation at 350 nm) spectra of dyad **1** (10 μ M) in MeCN.

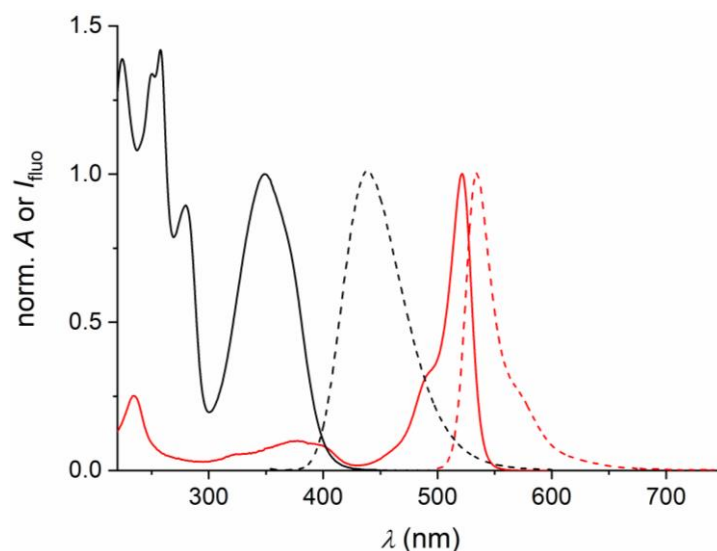


Fig. 3. UV/Vis absorption (solid line) and fluorescence (dashed line) spectra of model chromophores **2** (black) and **3** (red) (10 μ M each) in MeCN. The fluorescence spectra were obtained by excitation at 350 nm (**2**) or 495 nm (**3**).

Expectedly, only the Prodan-derived dye **2** and the corresponding part of dyad **1** showed significant solvatochromic effects of the absorption maxima, while the BODIPY chromophore exhibited the widely known polarity-independent spectral features. For dye **2** this is experimentally underpinned by a bathochromic red shift of about 30 nm upon changing from the less polar THF to the highly polar MeOH. Similar trends were observed for dyad **1**, with significant shifts of the longest wavelength absorption of the ICT chromophore ($\Delta\lambda = +15$ nm for MeOH *versus* THF) and practically solvent-independent absorption spectral features of the BODIPY moiety.

Table 1. Photophysical properties of **1–3** in various solvents.

	Solvent	λ_{abs} (nm) [ϵ ($\text{M}^{-1}\text{cm}^{-1}$)]	λ_{fluo} (nm)	Φ_{fluo}	τ_{fluo} (ns)
1	MeCN	350 [21700]	534 ^a	0.21 ^b	2.43 ^b
		522 [64000]		0.33 ^c	2.45 ^{c,d}
	THF	348 [19800]	537 ^a	0.53 ^b	5.00 ^b
		524 [74200]		0.77 ^c	5.14 ^c
	DMF	356 [22200]	538 ^a	0.20 ^b	2.53 ^b
		524 [66100]		0.34 ^c	2.62 ^{c,d}
	MeOH	363 [19900]	536	0.36 ^b	3.67 ^b
		523 [67300]		0.52 ^c	3.66 ^c
2	MeCN	349 [19100]	440	0.73	3.05
	THF	349 [19000]	422	0.55	2.58 ^d
	DMF	362 [19100]	450	0.62	3.22
	MeOH	378 [19700]	502	0.26	2.52 ^d
3	MeCN	521 [69600]	535	0.80	5.56
	THF	524 [69900]	538	0.82	3.92
	DMF	524 [77800]	537	0.77	5.14
	MeOH	522 [74000]	534	0.84	5.78

^a A very minor band corresponding to the antenna fluorescence was noted. THF: $\lambda_{\text{fluo}} = 423$ nm ($I_{537}/I_{423} = 250$); MeCN: $\lambda_{\text{fluo}} = 439$ nm ($I_{534}/I_{439} = 71$); DMF: $\lambda_{\text{fluo}} = 452$ nm ($I_{538}/I_{452} = 25$). ^b Values correspond to the excitation of the antenna chromophore at 350 nm. ^c Values correspond to the direct excitation of the BODIPY chromophore at 495 nm. ^d Biexponential decay; the average lifetime is given.

Furthermore, the UV/Vis absorption spectrum of the dyad resembles the sum of the two model chromophores and the high molar absorption coefficients of each moiety are also observed for the dyad (see Table 1). This corroborates the independence of the ground-state chromophores in **1**, which is an expected finding for an electronically isolating linker as the one employed herein.

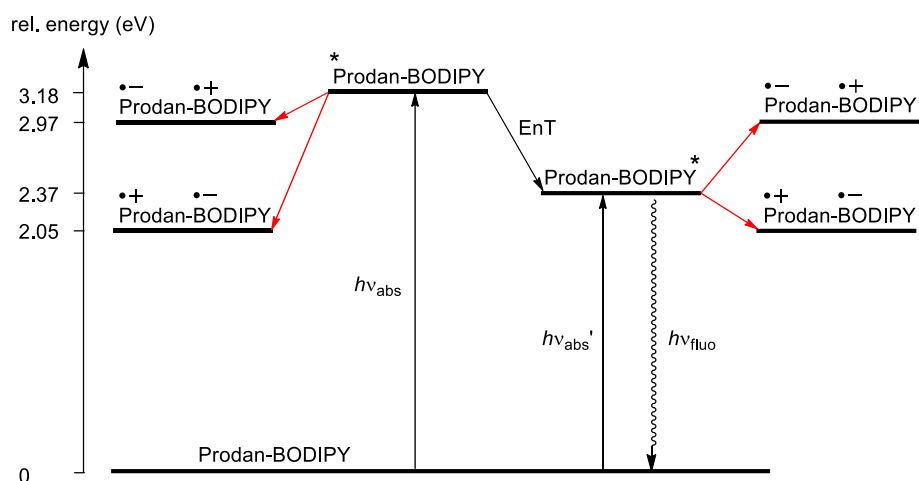


Fig. 4. Quantitative energy diagram depicting the hypothetical photoinduced electron transfer pathways (red arrows) in dyad **1**. Redox potentials (vs. SCE, in MeCN): E_{ox} (Prodan) = +0.74 V [63], E_{ox} (BODIPY) = +0.94 V [64], E_{red} (Prodan) = -2.09 V [65], E_{red} (BODIPY) = -1.37 V [64]. Excitation energies: E^* (Prodan) = 3.18 eV (calculated as average of energy of absorption and emission maximum; see ref. [66]), E^* (BODIPY) = 2.37 eV [64].

Regarding the fluorescence properties, the models **2** and **3** show again their characteristic behavior in solvents of different polarity: a bathochromic shift in polar solvents for the ICT chromophore **2** (up to 80 nm on going from THF to MeOH) and indifferent emission features of the BODIPY **3**. Strikingly, in the dyad **1** the fluorescence of the Prodan chromophore is hardly visible on its selective excitation (at 350 nm *ca.* 80% of the absorbed photons correspond to the Prodan-derived moiety), being in agreement with efficient quenching by excited state communication with the BODIPY chromophore (see Fig. 2). Instead practically only the BODIPY emission is observed, exhibiting the same spectral features as obtained for the direct excitation of the boron-containing fluorophore at 495 nm. This hints on an efficient quenching process involving excited-state EnT, where the ICT chromophore plays the role of the

energy donor and the BODIPY is the acceptor (see below). In comparison to model **3** the fluorescence quantum yields on direct excitation of the BODIPY in dyad **1** remain high in non-polar THF (0.77 for **1** *versus* 0.82 for **3**), but are considerably lower in polar solvents (*ca.* 0.3 for **1** *versus ca.* 0.8 for **3** in MeCN or DMF). Similar quenching effects are expressed by the decreased BODIPY fluorescence lifetimes of **1** in polar solvents. This fluorescence quenching is even more accentuated when exciting the Prodan-derived chromophore moiety instead, with BODIPY fluorescence quantum yields that are significantly smaller than measured for the direct BODIPY excitation in the dyad (see Table 1). The pronounced solvent dependence of these effects provides a hint that photoinduced electron transfer (PeT) pathways may be at their origin. This possibility was more carefully examined by taking into account the thermodynamic feasibility of the different PeT pathways (see Fig. 4). The driving force ΔG_{PeT} for the different PeT pathways was estimated by means of the Rehm-Weller equation [Eq. 1; E_{ox} – oxidation potential of electron donor, E_{red} – reduction potential of the electron acceptor, E^* – singlet excitation energy, C – Coulomb term (–0.06 eV in MeCN)] [67].

$$\Delta G_{\text{PeT}} = E_{\text{ox}} - E_{\text{red}} - E^* + C \quad (\text{Eq. 1})$$

The oxidative and reductive electron-transfer fluorescence quenching of the ICT chromophore are exergonic processes ($\Delta G_{\text{PeT}} = -1.13$ eV and -0.21 eV, respectively) and thereby energetically allowed. Along the same lines of thought, the reduced BODIPY fluorescence quantum yield obtained on direct excitation can be reasoned with an exergonic reductive PeT ($\Delta G_{\text{PeT}} = -0.32$ eV), while the oxidative pathway would be thermodynamically disfavored ($\Delta G_{\text{PeT}} = +0.60$ eV). Based on this analysis, it reasonable to assume that in polar solvents PeT competes with the EnT process and leads to the

observed quenching of BODIPY fluorescence (generated by direct excitation or by the EnT-sensitized pathway).

Table 2. Parameters of Förster resonance energy transfer in dyad **1** in various solvents.

	MeCN	THF	DMF	MeOH
$J \times 10^{-11}$ (cm ⁶ mol ⁻¹) ^a	4.0	1.8	5.9	16
R_0 (Å) ^b	35.3	28.6	35.2	37.7
$k_{\text{EnT}} \times 10^{11}$ (s ⁻¹) ^c	7.6	2.5	7.1	14
Φ_{EnT} ^d	1	1	1	1

^a Spectral overlap integral between donor emission and acceptor absorption spectra (see Eq. 2). ^b Critical FRET radius (see Eq. 3), with $\kappa^2 = 0.2461$. ^c Energy-transfer rate constant (see Eq. 4). ^d Energy-transfer quantum yield (see Eq. 5).

Albeit the above discussed PeT pathways interfere in the excited-state behavior of **1** (especially in polar solvents), it is clear from the abovementioned observations that excitation of the Prodan-derived moiety leads to the observation of BODIPY fluorescence emission. This can be only reasonably explained with EnT communication between both chromophores. With the objective to shine more light on this process, a more detailed analysis according to the Förster theory of resonance energy transfer (FRET) was undertaken [68]. The key data, such as the spectral overlap integral J , the critical transfer radius R_0 , the quantum yield Φ_{EnT} , and the rate constant k_{EnT} of energy transfer are summarized in Table 2; see equations 2–5.

$$J = \int_0^{\infty} [F_D(\nu)\varepsilon_A(\nu)d\nu]/\nu^4 \quad (\text{Eq. 2})$$

$$R_0^6 = 8.8 \times 10^{-28} \Phi_D \kappa^2 n^{-4} J \quad (\text{Eq. 3})$$

$$\Phi_{\text{EnT}} = 1/[1 + (R/R_0)^6] \quad (\text{Eq. 4})$$

$$k_{\text{EnT}} = (1/\tau_{\text{D}}) \times (R_0/R)^6 \quad (\text{Eq. 5})$$

The actual center-to-center distance of the donor and acceptor units, obtained from molecular dynamics simulations (see details in the Supplementary Material), accounting for the conformational flexibility of the linker, is significantly smaller than the critical transfer radius, e.g., 9.8 Å *versus* 35.3 Å in MeCN. This leads to the prediction of practically quantitative inherent EnT (Φ_{EnT} *ca.* 1), independent on the polarity of the solvent. In practice, PeT (see above) competes with this process, thereby lowering the quantum yields of sensitized BODIPY fluorescence in polar solvents. However, in all solvents the EnT is evidently operative and in essence the relatively small Stokes shift of the BODIPY chromophore itself (*ca.* 10 nm) is considerably enlarged to *ca.* 170–180 nm in the dyad, referring to an apparent Stokes shift.

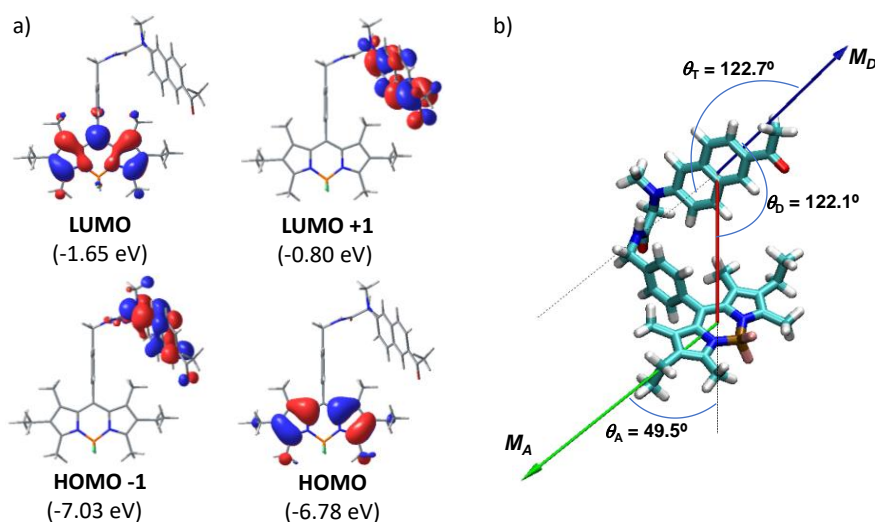


Fig. 5. (a) Contour plots of relevant frontier orbitals of dye **1**, and (b) Optimized ground-state geometry showing data for calculation of the orientation factor κ^2 corresponding to dyad **1**.

Table 3. Main absorption and emission transitions of dyad **1** and the corresponding model compounds **2** and **3**.

Dyad 1	2	3
Absorption		
$S_1 \leftarrow S_0$ ($f = 0.680$) ^a LUMO \leftarrow HOMO (98%) $E = 2.92$ (2.38) ^b eV $S_3 \leftarrow S_0$ ($f = 0.472$) ^a LUMO+1 \leftarrow HOMO-1 (90%) $E = 3.70$ (3.54) ^b eV	$S_1 \leftarrow S_0$ ($f = 0.487$) ^a LUMO \leftarrow HOMO (90%) $E = 3.66$ (3.55) ^b eV	$S_1 \leftarrow S_0$ ($f = 0.691$) ^a LUMO \leftarrow HOMO (98%) $E = 2.92$ (2.38) ^b eV
Emission		
$S_1 \rightarrow S_0$ ($f = 0.837$) ^a LUMO \rightarrow HOMO (98%) $E = 2.80$ (2.32) ^b eV $S_3 \rightarrow S_0$ ($f = 0.728$) ^a LUMO+1 \rightarrow HOMO-1 (93%) $E = 2.98$ (2.82) ^b eV	$S_1 \rightarrow S_0$ ($f = 0.808$) ^a LUMO \rightarrow HOMO (96%) $E = 2.99$ (2.82) ^b eV	$S_1 \rightarrow S_0$ ($f = 0.840$) ^a LUMO \rightarrow HOMO (98%) $E = 2.80$ (2.32) ^b eV

^a Oscillator strength f . ^b In parenthesis the experimental transition energy, based on the absorption and emission maxima, is given.

3.3. Theoretical calculations

The electronic independence of the two chromophores in dyad **1** was confirmed by time-dependent density functional theory (TD-DFT) calculations at the CAM-B3LYP/6-31+G(d) level of theory. The solvent (MeCN) was taken into account by means of the polarizable continuum model (PCM). The calculations confirmed orthogonal chromophore-specific transitions with no sign of mixing: $S_3 \leftarrow S_0$, dominated by LUMO+1 \leftarrow HOMO-1 (90% weighted contribution) for the Prodan-derived moiety and $S_1 \leftarrow S_0$, dominated by LUMO \leftarrow HOMO (98% weighted contribution) for the BODIPY moiety (see Fig. 5). The corresponding frontier molecular orbitals (FMO) are localized on the respective chromophoric unit, which clearly underpins the electronic insulation by the non-conjugative linker. The calculated excitation and emission energies of **1** express the experimentally verified trends correctly (see Table 3). Further,

a comparison with the calculated data for the models **2** and **3** underpins the electronic orthogonality of the two chromophores in dyad **1** (Supplementary Material – Fig. S33–S34).

3.4. Two-photon excitation (TPE)

The electronic constellation of the Prodan-derived antenna chromophore, being a push-pull architecture, opens the possibility for exploring two-photon absorption (TPA) processes and their implication in the EnT-sensitized generation of BODIPY fluorescence. As exemplified for model compound **3**, the plain BODIPY chromophore is a rather inefficient two-photon absorber (<10 GM at wavelengths between 740 and 1040 nm). However, by the target-oriented design of electronic properties an improvement of the TPA properties of BODIPY dyes can be achieved [3,69,70]. The herein followed strategy is different. It relies on the two-photon-induced population of the first excited singlet state of the antenna chromophore that engages subsequently in the same EnT process as observed under one-photo-excitation (OPE) conditions (see above). The TPA action spectrum of dyad **1** in DMF was obtained by the comparative method using rhodamine B as reference (see Materials and methods). The data are shown in Fig. 6. Two spectrally well differentiated maxima at 700 nm and 970 nm are observed. In agreement with the spectra that are known for the two separate moieties of the dyad (see Fig. 6) these bands are assigned to the Prodan-derived antenna chromophore and the BODIPY moiety, respectively. Having in mind that the BODIPY fluorescence quantum yield varies in dependence on the fact whether it was generated by a sensitization mechanism or direct excitation (see above), the absolute two-photon cross section σ_2 of the antenna part is calculated as 293 GM at 700 nm and the one of the BODIPY moiety as 32 GM at 970 nm (1 GM = 10^{-50} cm⁴ s photon⁻¹). The BODIPY

emission spectra that are obtained on excitation in either TPA band (at 700 nm or 1000 nm) are identical (see Fig. S37 in the Supplementary Material), showing the typical narrow-shaped profile and a maximum at *ca.* 540 nm. Akin to the observations made for OPE (see above), for the excitation of the Prodan-derived moiety at 700 nm a very minor emission with a maximum around 455 nm was seen. These joint observations underpin our notion of the operation of the same EnT process that was characterized in detail for the OPE conditions (see above). In fact, the attachment of the Prodan-derived antenna does not only increase the apparent Stokes shift (see above) of the BODIPY dye, but also improves indirectly its TPA properties by means of the antenna effect.

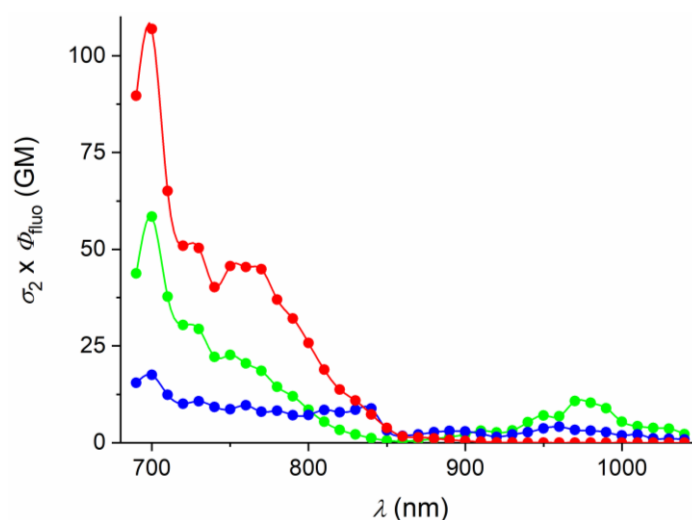


Fig. 6. Two-photon absorption spectra for bichromophoric dyad **1** (green), and models **2** (red) and **3** (blue) in DMF (10 μM each).

3.5. Cellular imaging with dyad **1**

With view on potential bioimaging applications dyad **1** was investigated regarding the operation of the two-photon-induced EnT in the environment of the N13 microglial cells. The incubation of dyad **1** with N13 microglial cells, following a standard protocol (see Materials and methods), and subsequent fluorescence imaging on excitation at 750

nm (absorbed predominantly by the Prodan-derived moiety) yielded well resolved images that evidenced the successful cellular uptake of the probe and its accumulation in the cytoplasm (see Fig. 7 and the Supplementary Material for bright-field images – Fig. S38). The brightness of the dyad is sufficiently high to lower the concentration of the dye in the incubation broth down to 1 μM .

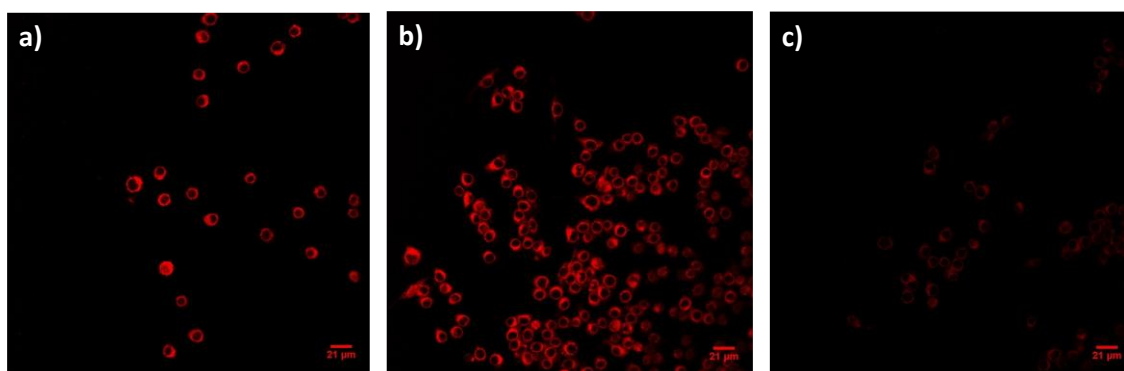


Fig. 7. Two-photon confocal microscopy images of N13 microglial cells incubated with dyad **1** in different concentrations; excitation at 750 nm. Concentrations of **1**: 100 μM (a), 10 μM (b), 1 μM (c).

The *in cellulo* detected fluorescence spectrum revealed the typical narrow emission band of the BODIPY with a maximum at about 540 nm (see Fig. 8). Furthermore, the two-photon excitation spectrum showed a strong band with a maximum at *ca.* 720 nm and a smaller shoulder at *ca.* 790 nm, assigned to the Prodan-derived chromophore (see Fig. 8). The BODIPY chromophore was represented by a weak band at *ca.* 990 nm. These features coincide with the observations in solution. Hence, it can be safely concluded that the photophysical key features of the bichromophoric dyad **1**, namely TPA and the EnT communication, persist under the experimental conditions of fluorescence imaging in the model cell line.

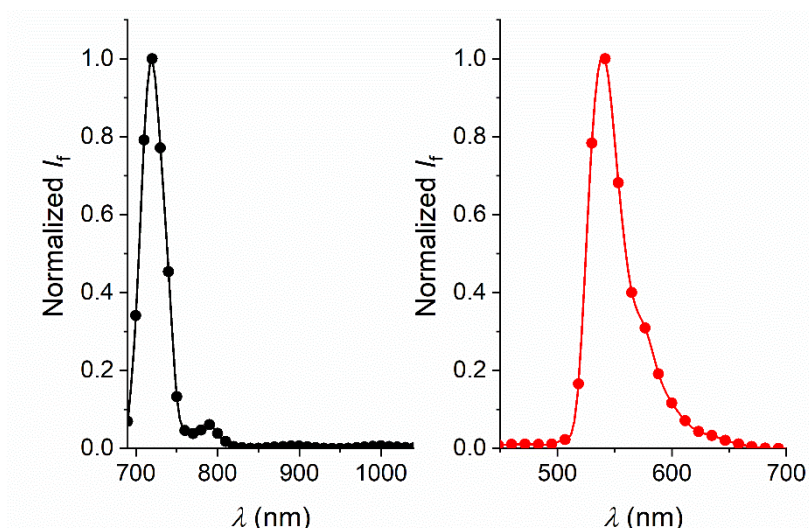


Fig. 8. Left: Two-photon excitation spectrum of **1**, incubated with N13 microglial cells (monitoring of emission between 542 and 572 nm). Right: Corresponding emission spectrum, excitation at 750 nm.

Conclusions

The combination of a two-photon-absorbing antenna Prodan dye with a BODIPY fluorophore permitted the demonstration of the EnT-sensitized fluorescence emission of the latter on excitation of the antenna with near-infrared photons. The EnT is predicted to be highly efficient. However, in polar solvents the competition with PeT processes adds some complexity to the excited state characteristics. The EnT-sensitization is still operative *in cellulo*, which permits the use of the dyad in bioimaging.

CRedit authorship contribution statement

Antonio Dominguez-Alfaro: Investigation. **Vânia Pais:** Investigation, Validation.

David B. Guzmán-Ríos: Investigation, Writing - Review & Editing. **Daniel Collado:**

Investigation. **Francisco Nájera:** Investigation, Writing - Review & Editing. **Ezequiel**

Pérez-Inestrosa: Supervision, Writing - Review & Editing. **Uwe Pischel:**

Conceptualization, Supervision, Writing - Original Draft.

Declaration of competing interest

The authors declare that they have no competing financial interests or personal relationships that could have influenced the work reported in this paper.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/...>

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