

Band Gap Narrowing vs. Formation of Electronic States in the Gap in N-TiO₂ Thin Films

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3 **Band gap narrowing vs. formation of electronic states in the gap in N-TiO₂ thin**
4 **films**
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28 **Abstract**
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31 N-containing TiO₂ thin films with different amounts of nitrogen have been prepared by
32 plasma enhanced chemical vapor deposition (PECVD) by using different titanium
33 precursors without (titanium isopropoxide, TTIP) and with (tetrakis diethyl amino
34 titanium, TDEAT and tetrakis dimethyl amino titanium, TDMAT) nitrogen in their
35 structures and different N₂/O₂ ratios as plasma gas. For low/high content of nitrogen, Ti-
36 NO and/or Ti-N like species have been detected in the films by X-ray photoelectron
37 spectroscopy (XPS). Their optical behavior is characterized by a red shift of their
38 absorption edge when Ti-N species are a majority and by an unmodified edge with
39 localized absorption states in the gap when only Ti-NO like species are present in the
40 film. The experimental results have been interpreted by calculating the density of states
41 of model systems consisting of a 2x2x3 repetition of the anatase unit cell. This basic
42 structure incorporates nitrogen defects in either substitutional or interstitial lattice
43 positions that are considered equivalent to the Ti-N and Ti-NO like species detected by
44 XPS. To simulate the effect of, respectively, a low or a high concentration of nitrogen,
45 calculations have been carried out by placing two nitrogen defects either in separated or
46 in nearby positions of the anatase structure. The computational analysis reveals that the
47 defects have different stabilization energies and confirm that an edge shift of the
48 valence band is induced by the substitutional nitrogen centers, as observed when a high
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3 concentration of Ti-N species becomes incorporated into the films. In agreement with
4 the experimental results, when only Ti-NO like species are detected by XPS no band
5 gap narrowing is obtained by the calculations that predict the appearance of localized
6 electronic states in the gap. The fact that only these latter films present water wetting
7 angle photo-activity when irradiated with visible light supports that the presence of Ti-
8 NO like species is a required condition for visible light photo-activity.
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16 **Keywords:** N-modified TiO₂; plasma deposition; optical properties; visible photo-
17 activity; electronic structure; N species
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23 **1. Introduction**

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26 N-modified TiO₂ materials have been widely studied during the last years because of
27 the possibility to make them photo-catalytically active under visible light illumination.
28 Since the seminal work of Asahi et al,¹ many papers have appeared in the literature
29 trying to unravel the reasons that sustain this behavior.²⁻¹⁰ Of particular relevance in
30 this regard are the recent works by Di Valentin et al¹¹⁻¹⁵ where, by combining theory
31 and experiments mainly by EPR, they show that the species responsible for the visible
32 photo-activity of N-containing TiO₂ consist of N atoms located in interstitial lattice
33 positions where it bonds to Ti atoms and interacts with a nearby O ion of the lattice
34 (i.e., a kind of Ti-NO species). This model contrasts with more traditional views linking
35 the visible photo-activity of N-doped TiO₂ with the formation of substitutional nitrogen
36 species (i.e., a kind of Ti-N species).^{1, 5,16} Another related aspect that has generated
37 much discussion within this topic is the influence of the incorporated nitrogen on the
38 optical properties of the materials, namely in the changes in light absorption due to
39 possible modifications of the band gap of the titanium oxide.¹⁷⁻²¹ The appearance of
40 absorption features in the visible region of the spectra of N-containing TiO₂ is generally
41 taken as a required but not sufficient condition for visible photo-activity.²² However, the
42 physical causes of this visible absorption are also controversial, with some authors
43 supporting that the effect of the nitrogen species is to shift the band gap of the titanium
44 oxide towards the visible¹⁷⁻²⁰ and others claiming that these species do not modify the
45 band gap of the oxide but induce some electronic states in the band gap region that give
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3 rise to localized absorption features close to the absorption edge of the titanium oxide.¹¹⁻
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8 In a recent work on N-doped TiO₂ thin films prepared by plasma and evaporation
9 methods, we have shown that while Ti-NO species are able to induce some visible light
10 decrease of the wetting angle of the surface,²² Ti-N species do not induce any kind of
11 visible light activation. However, for the samples to present photo-catalytic activity
12 towards the photo-degradation of dyes it was required that besides Ti-NO species, the
13 films were crystalline with a mixture of anatase and rutile. In this previous investigation
14 we already noted that the absorption spectra of films with a high content of nitrogen
15 depicted some absorption features in the visible, although no correlation could be
16 established between the nitrogen content and the characteristics of the UV-vis
17 absorption spectra.
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22 As a continuation of this previous work, herein we aim primarily at determining
23 whether each kind of nitrogen species (i.e. Ti-NO or Ti-N) has a specific effect on the
24 physical causes (i.e., gap narrowing or development of localized states in the gap)
25 determining the changes in the absorption spectra of N containing titanium oxide thin
26 films (in the following named in general as N-TiO₂ although changes in the Ti/O ratio
27 may occur as a consequence of the development of oxygen vacancies whose formation
28 is related with the incorporation of Ti-N species of nitrogen²⁵). We approach this
29 problem firstly by investigating experimentally a series of N-TiO₂ thin films where we
30 have systematically changed the type and concentration of nitrogen species. To get such
31 an ample set of samples with different concentrations of nitrogen species we have
32 employed the plasma enhanced chemical vapor deposition (PECVD) procedure and
33 modified systematically the plasma conditions and other working parameters like the
34 type of precursor and temperature of the substrate. Such an experimental approach has
35 hitherto not been widely used for the preparation of photoactive N-TiO₂ thin films,
36 although the preparation of photo-active nanoparticles by atmospheric plasma
37 deposition has been reported.²⁶ The most common titanium precursor used for the
38 PECVD synthesis of TiO₂ thin films is titanium isopropoxide (TTIP),^{23, 27-30} although
39 tetrakis dimethyl amino titanium (TDMAT) and tetrakis(diethylamido) titanium
40 (TDEAT)³¹⁻³³ are other two precursors which are commonly used in methods of
41 preparation of optical thin films. Therefore, another innovative approach of our work is
42 the use of three titanium precursors (TTIP, TDEAT and TDMAT) with the idea that the
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3 presence of direct Ti-N bonds in the structure of two of them and, very likely, their
4 different reactivity, permit to tune the state of nitrogen in the films and to control film
5 properties such as UV-vis absorption spectra and photo-activity.
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10 To account for the different experimental findings and particularly for the changes in
11 the absorption spectra depending on the type and concentration of nitrogen species in
12 the prepared N-TiO₂ thin films, we have theoretically modeled the electronic structure
13 of these materials by calculating the density of states of N-modified anatase for three
14 different types of N defects incorporated in its structure. Modeling the Ti-NO (and Ti-
15 ON as a possible alternative) defect states is done by assuming that N occupies an
16 interstitial position and interacts with both oxygen and titanium ions of the lattice. This
17 species, with nitrogen bearing a formal charge of -1 (i.e., NO³⁻ like species), could be
18 described as a (NO)_O^t defect centre according to the Kröger-Vink notation of defect
19 states in oxides. To simulate the Ti-N species, nitrogen replaces substitutionally an
20 oxygen of the lattice and bonds directly to titanium in a kind of nitride species (i.e. N³⁻
21 species or N_O^t according to the Kröger-Vink notation). Similar calculations simulating
22 the effect of Ti-N and Ti-NO like species can be found in the works by Di Valentin et
23 al.^{11,15} and Graciani et al.^{25,34} Herein, besides studying the stability of Ti-ON like
24 species, we have simulated the effect on the optical properties of N-TiO₂ materials of
25 the concentration of Ti-N or Ti-NO species, an aspect that had not been addressed in
26 these previous theoretical works. For this purpose, calculations have been carried out by
27 placing two defect states quite separated (i.e., simulating a low concentration of
28 nitrogen) or in nearby positions of the structure (i.e., simulating a high concentration of
29 nitrogen), in this latter case making it possible that the two defect states interact
30 between them as it would be the case in samples with a high concentration of nitrogen
31 species. The density of state calculations (DOS) and particularly the analysis of the
32 partial density of state plots of N, Ti and O have confirmed the experimental findings
33 and show that the incorporation of a high concentration of Ti-N species induces a red
34 shift in the absorption gap of the material, while Ti-NO species do not significantly
35 modify the magnitude of the gap between the valence and conduction bands, but induce
36 the formation of new localized states in the band gap zone of the oxide.
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58 Finally, all these experimental and theoretical results are critically discussed in relation
59 with the actual photo-activity of N-TiO₂ materials by analyzing the super-hydrophilic
60 conversion of the surface state of the prepared films under visible light irradiation.

2. Experimental

2.1 *Thin film preparation*

TiO₂ and N-TiO₂ thin films have been prepared by PECVD in a plasma reactor with a remote configuration. The system, supplied with a microwave plasma source (SLAN, from Plasma Consult, GMBh, Germany) has been described in previous works.^{27, 28, 35} It consists of an external 2.45 GHz microwave electron cyclotron resonance (MW-ECR) plasma source coupled to the reaction chamber and separated from it by a grid to avoid the microwave heating of the substrates. Under normal conditions of operation, the grid confines the plasma out of the reaction chamber (remote plasma conditions) where the substrate and the precursor dispenser are located. Titanium tetra-isopropoxide (TTIP), Tetrakis(diethylamido)titanium(IV) (TDEAT) and Tetrakis(dimethylamido)titanium(IV) (TDMAT) were used as titanium precursors. In Figure 1 we show a schematic representation of the structure of these three precursors, showing that four direct Ti-N bonds exist in TDEAT and TDMAT. It is likely that some of them can be preserved during the plasma processing of the N-TiO₂ thin films.

The plasma source was operated with a power of 400 W with either pure O₂ or mixtures N₂+O₂ at different proportions of nitrogen from 0 % to 97%. The synthesis of the films was carried out either at 298 K or 523 K, although most reported data corresponds to this latter temperature. Controlled dosing of the precursors was achieved by placing them in a stainless steel recipient heated to 305 K, while oxygen or a N₂/O₂ mixture was bubbled through it. Both the bubbling line and the shower-type dispenser used to dose the precursor into the chamber were heated to 373K to prevent any condensation in the tube walls. The total pressure during deposition was 4x10⁻³ Torr (normal operation conditions). A deposition rate of approximately 2.5 nm min⁻¹ was estimated by means of a quartz crystal monitor for the samples grown at room temperature.

The different plasmas used for deposition of the films were characterized by Optical Emission Spectroscopy (OES). The emission spectrum was analyzed using a 0.5 m CVI/Digikrom DK480 monochromator (CVI Laser Corporation, Albuquerque, NM) with a 1200 grooves/mm grating, spectra resolution of 0.2 nm, and spectral sensitivity ranging from 200 to 900 nm (Hamamatsu R928). The light is collected by a fiber optic,

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3 which is fixed to the monochromator and positioned 4 cm downstream from the source
4 center. The light was collected through a hole in the air cooling ring. It should be noted
5 that the spectroscopic intensities are integrated over the line of sight, thus limiting the
6 spatial resolution.
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10 All the films were deposited simultaneously on a silicon wafer and on quartz plates. The
11 thickness of the prepared samples was estimated by measuring the mass thickness of the
12 films by X-ray fluorescence (XRF) and Rutherford Back Scattering (RBS) or directly
13 from cross section views in the electron microscope (note that for a given sample these
14 two values do not coincide except for completely compact thin films). Most
15 experiments were carried out with films with a thickness of some hundreds of
16 nanometers.
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23 *2.2 Methods of characterization*

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26 The optical properties of the samples were primarily determined by UV-vis absorption
27 spectroscopy (Perkin-Elmer Lambda 12 Spectrometer) for samples prepared on fused
28 silica. The ellipsometric characterization of the N-TiO₂ films was carried out in a J.A.
29 Woollam VASE (variable angle spectroscopic ellipsometry) spectroscopic ellipsometer.
30 Values of Ψ and Δ were obtained over the spectral range comprised between 300 and
31 1000 nm, at 5nm resolution. To check the consistency of the data, they were collected at
32 three angles of incidence: 65°, 70°, and 75°. Optical modeling and parameters fitting
33 were carried out with the WASE32© program (J.A. Woollam Co.,Inc.). To model the
34 ellipsometric spectra we fixed the thickness (thickness values were obtained by
35 transversal SEM images) and applied non-ideal model options such as film thickness
36 non-uniformity and angular spread of the beam entering the detector. Some fitting
37 parameters were coupled among them and their range limited within defined values.
38 Non partial polarization or monochromator bandwidth effects were considered. Quality
39 assessment of the fit data was performed by using the mean-squared error (MSE) value.
40 A small MSE implies that the assumed model was appropriate. In our case, the MSE
41 was below 5 units in the absorbing region and below 2 units in the transparent region.
42 The used optical model consists of two layers deposited on the Si (1 mm) substrate: an
43 external region and an EMA (effective medium approximation) layer with a harmonic
44 oscillator. In this layer a Cauchy dispersion (with Urbach absorption)) and a classic
45 harmonic oscillator model were implemented in a single material layer. The Cauchy
46 dispersion equation was used in the transparent region ($\lambda > 700\text{nm}$). No dispersion of
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3 refractive index and in the extinction coefficient were considered for $700 < \lambda < 1000 \text{ nm}$.
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5 For $\lambda < 700 \text{ nm}$ one or two harmonic oscillators were added to the Urbach absorption of
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7 the Cauchy model. Finally, a certain surface roughness consisting of 50% of film
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9 material and 50% of voids was assumed to describe the external layer.

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11 XPS spectra of the films were recorded on an ESCALAB 210 spectrometer working
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13 under energy transmission constant conditions. The Mg K α line was used for excitation
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15 of the spectra. They were calibrated in binding energy (BE) by referencing the different
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17 peaks to the C1s signal due to contamination taken at 284.6 eV. Quantification was
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19 done by calculating the area of the peaks and by correcting then with the sensitivity
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21 factor of each element/electronic level. To check the effect of surface impurities, the
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23 films were subjected to a gentle sputtering with Ar⁺ ions of 2.5 keV. A current density
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25 of about 10 $\mu\text{A cm}^{-2}$ for a sputtering time of 2 min was used for these treatments. No
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27 significant changes in peak shape or in element ratio (except for the removal of most of
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29 the contaminating carbon) were found after this cleaning treatment. Fitting analysis of
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31 the N1s peak was carried out by using elemental bands of gaussian/lorentzian shape
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33 after background subtraction of the spectra with a Shirley-type curve.

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35 SEM cross section and normal images were measured in a Hitachi S5200 field emission
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37 microscope for thin films grown on a silicon wafer.

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39 Structural characterization of the thin films was done by X-ray Diffraction in a Siemens
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41 D5000 diffractometer working in the Bragg-Brentano configuration and using the Cu K α
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43 line as excitation source.

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45 Measurement of the surface electrical conductivity of the samples was carried out with
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47 the typical four point probe test. A Keithley 617 Electrometer and a Hewlett-Packard
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49 34401 A Voltammeter were used for the measurements. These consisted of applying a
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51 voltage ranging between -0.25 and 0.25 V to the two external probes and the
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53 measurement of the current flowing between the two internal probes.

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55 Measurement of water contact angles was carried out by the Young method by dosing
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57 small droplets of deionized and bidistilled water on the surface of the illuminated
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59 samples. In the experiments where the contact angle variation was determined as a
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61 function of the illumination time, a metal foil acting as a shutter was used to close and
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63 open the lamp output. All wetting angle measurements within a given experiment were
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65 taken after illumination for successive periods of time. Therefore, the time scale in the
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67 plots refers to the accumulative illumination of the samples. The maximum uncertainty

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3 in the determination of the water contact angle is about 10° depending on the sample
4 position. In the course of this investigation it was noticed that the “as-prepared” thin
5 films were more hydrophilic than the same samples some time after their preparation.
6 Therefore, the reported results correspond to samples that were stored in desiccators for
7 at least two months before testing their photo-activity. Illumination of the samples was
8 carried out with a Xe discharge lamp with photon intensity at the position of the
9 samples of 2 Wcm^{-2} for the complete spectrum. For simplicity we will refer this
10 situation in the text and figures as UV illumination. Other experiments consisted of the
11 illumination with the same lamp by placing an UV filter (i.e., $\lambda > 400 \text{ nm}$) between the
12 lamp and the sample. The light intensity at the sample position was then 1.6 Wcm^{-2} . In
13 all cases an infrared filter (i.e., a water bath) was kept between the lamp and the samples
14 to prevent any possible heating by the infrared radiation.
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25 *2.3 Computational details*

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28 All the calculations were performed using the Vienna Ab-initio Simulation Package
29 (VASP) 4.6 code,³⁶⁻³⁸. Inner core electrons were described by the projector-augmented-
30 wave method⁴¹ (PAW) and a 500 eV cut-off energy. The Ti (3s, 3p, 3d, 4s), O (2s, 2p)
31 and N (2s, 2p) electrons were treated as valence states, while the remaining electrons
32 were kept frozen as core states. The optimisation of the structures was performed via a
33 conjugated gradient technique, which utilises the total energy and the Hellmann-
34 Feynman forces on the atoms (and stresses on the unit cell). The Brillouin-zone
35 integrations were performed using Monkhorst-Pack grids,⁴² using a (2x2x1) mesh, with
36 which convergence in energy was achieved. The simulation cell is formed by a 2x2x3
37 repetition of the anatase unit cell, in which there are 72 Ti atoms and 144 O atoms. We
38 have employed a LDA+U approach, as described by Dudarev et al,⁴³⁻⁴⁵ with a value of
39 U of 4.5 eV. Calzado et al.⁴⁶ found that this value of U yields a correct description of
40 the gap states in periodic LDA+U calculations. All calculations have been carried out in
41 the Finis Terrae supercomputer, in Santiago de Compostela, Spain.
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3. Results

3.1 *State and concentration of nitrogen in N-TiO₂ thin films*

The amount of nitrogen incorporated in the film varied by changing the precursor and the plasma process parameters. Adjusting the plasma gas composition by using different O₂+N₂ mixtures as plasma gas was very important for this purpose. The OES spectra of the plasmas of these mixtures for N₂ concentration higher than 80% were characterized by a series of peaks and bands attributed to N₂* and N₂⁺ species (see supporting information S1).^{47,48} These nitrogen bands substitute those attributed to O₂* and O* that characterize the spectra recorded in our experimental system when pure oxygen is used as plasma gas.²⁷ A summary of the main species and spectral features detected as a function of the O₂/N₂ ratio are reported in the supporting information S1 and in a previous contribution of our group.⁴⁹ From this analysis it is important to realize that although nitrogen excited species detected by OES are majority in the N₂/O₂ plasmas, this does not necessarily preclude the formation of activated oxygen species, but just indicates that the average life time of these latter species are short as compared with that of the nitrogen species. The high reactivity of the oxygen species accounts for the formation of a titanium oxide film and enables the release of the carbon and hydrogen atoms of the precursors in the form of volatile oxidized compounds. Simultaneously, the presence of active species of nitrogen in the plasma provides a mechanism for the incorporation of nitrogen in the film by the reaction with the partially decomposed precursor molecules in the plasma phase or on the surface of the growing film. We will show later that the amount and type of incorporated nitrogen depends on the ratio between nitrogen and oxygen in the plasma gas.

Plasma deposition of TiO₂ thin films has been typically carried out by using TTIP as titanium precursor. The use of other precursors like TDMAT and TDEAT offers the possibility to incorporate some nitrogen within the structure even if no molecular nitrogen is added to the plasma gas, provided that some of the Ti-N bonds existing in the precursor molecules are preserved in the films. To verify this possibility we have analyzed by XPS a series of films prepared with a plasma of pure oxygen but using TTIP, TDMAT and TDEAT as precursors and 523 and 298 K as substrates temperatures. It was found that small amounts of nitrogen became incorporated in the films prepared with TDMAT and TDEAT precursors (see supporting information S2).

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3 The incorporation was more favorable at 298 K. For all samples the O1s and Ti2p
4 spectra were similar to those amply reported in the literature for TiO₂.⁵⁰ The N/Ti ratios
5 deduced from the quantitative analysis of the spectra of the N-TiO₂ films prepared at
6 523 K and O₂ as the sole plasma gas were 0.037, 0.017 and 0.0 for the TDEAT,
7 TDMAT and the TTIP precursors, respectively. The N1s spectra of the two former films
8 were characterized by a broad peak at around 400 eV (see supporting information S2)
9 which is commonly attributed to interstitial nitrogen species bonded to titanium and
10 oxygen (i.e., a kind of Ti-N-O species^{12, 22}). Similar spectra were obtained when the
11 deposition was carried out with a substrate temperature of 298 K, although the intensity
12 of the N1s signal was approximately twice higher in this case. It is also important to
13 mention that the width of this N1s peak was also higher at 523 K (i.e., a HWHM of 2.8
14 eV) than at 298 K (i.e., a HWHM of 1.8 eV). This difference suggests that the nitrogen
15 signal at 523 K is a convolution of different species of nitrogen. N1s spectra with more
16 than one component at around 400 eV have been reported by several authors.⁵¹⁻⁵³ Their
17 development in our case would indicate that at 523 K there are nitrogen species with
18 slightly different geometries and bonding arrangements.

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33 The concentration and type of nitrogen species in the N-TiO₂ thin films changed when
34 mixtures O₂+N₂ were used as plasma gas. Under these conditions the main parameter
35 controlling the incorporation of nitrogen into the films was the concentration of nitrogen
36 in the gas mixture. Figure 2 shows a series of N1s spectra taken for films prepared at
37 523 K by using TDMAT as titanium precursor and increasing concentrations of N₂ in
38 the plasma gas. Similar spectra were obtained for the other precursors. It is found that
39 for percentages of N₂ higher than 85%, a new nitrogen species at around 396 eV appears
40 in the spectrum. This species becomes dominant for higher percentages of nitrogen in
41 the plasma gas. Similar results were obtained for the two other precursors, thus
42 sustaining that the plasma gas composition and not the type of precursor is the main
43 variable controlling the amount of nitrogen incorporated into these films. This new form
44 of nitrogen has been attributed on the basis of XPS studies to nitride species in
45 substitutional positions of the titanium oxide lattice (i.e., a Ti-N species^{1, 7, 10 53-55}). The
46 combined XPS, NEXAFS and Rietveld analysis made recently by H. Chen et al⁵⁵ for
47 ammonia treated titania have quite unequivocally supported this attribution. It is also
48 worthy of note that for most investigated films the shape of the O1s and Ti2p spectra is
49 similar to that reported for TiO₂, thus indicating that the upmost surface layers of the
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3 films have become oxidized by exposure to the air. Nevertheless, the shape of the Ti2p
4 spectrum of the films prepared with a 97% of nitrogen in the plasma gas slightly
5 changed with the development of a shoulder at lower binding energy, evidencing that
6 some Ti^{n+} species ($n < 4$) remain in the surface of this sample even after its handling in
7 air.
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12 A deeper insight into these spectra can be obtained by fitting. Fig. 3 shows such an
13 analysis for some selected case examples taken from Fig. 2 and others corresponding to
14 thin films prepared with TTIP as precursor. The two sets of samples were synthesized
15 with relative N_2 contents in the plasma gas of 0, 90 and 97%. In the case of the
16 TDMAT samples, the spectra can be properly fitted with four bands located at 396.1,
17 399.3, 400.7 and 401.7 eV, the three latter accounting for the broad peak at around 400
18 eV. Meanwhile, a reasonable fitting could be obtained for the spectra of the TTIP
19 samples by using just three bands at 396.1, 399.3 and 400.7 eV. A variety of peaks with
20 N1s BEs comprised between 399 and 402 eV were already found in the early work of
21 Saha et al.⁵⁶ dealing with the thermal oxidation of titanium nitride. Because of the
22 similarities of the preparation techniques, it is also worth mentioning a recent
23 publication by Chen et al.²⁶ where they investigated by XPS the type of nitrogen species
24 incorporated into N-TiO₂ powder samples prepared by an atmospheric plasma
25 procedure and posterior annealing. These authors found different components in the
26 N1s photoelectron spectra, which they attributed to Ti-N (395.8-397.8 eV), Ti-NO
27 (398.8-401.2 eV) and Ti-NO₂ (402.0-403.0 eV) species, in good agreement with our
28 own results. Another relevant result of that work was the verification that the latter
29 species was inactive in inducing any visible photo-catalytic activity. Similarly, Dunnill
30 et al.,⁵⁷ using a thermal chemical vapor deposition process at 500 °C for the preparation
31 of N-TiO₂ thin films, found that in the majority of cases a peak at 400 eV was the only
32 species detected and that these samples presented photocatalytic activity in the visible.
33 Based on these and many other recent studies,^{1, 7, 10, 12, 22, 51-55} we attribute the different
34 features to Ti-N (i.e., band at 396.3 eV assigned to nitrogen triple bonded to titanium)
35 and to nitrogen in a titanium oxinitride local environment (399.3 and 400.7 eV), where
36 nitrogen simultaneously bonds to oxygen and to titanium in a defective lattice site (i.e.,
37 in a kind of Ti-N-O or Ti-O-N local structure). Meanwhile, the fourth band at 401.7
38 eV, observed in samples prepared with the TDMAT and TDEAT precursors, must be
39 attributed to N bonded to a less electropositive environment or, alternatively, to a large
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3 number of oxygen atoms. Therefore, we attribute tentatively this band to Ti-NO₂
4 species (i.e. a nitrogen atom bonding simultaneously to one titanium and two oxide ions
5 of the lattice ²⁶). A similar attribution has been made recently by Asahi et al⁵³ based on
6 first principles calculations for a peak with a BE higher than that of the band used here
7 for fitting. For simplicity, in the rest of the paper we will refer to Ti-N and Ti-NO
8 species (this latter including the species that we have designed as Ti-ON), even if
9 different environments, and therefore BEs, can be expected for the different interstitial
10 species of nitrogen. Theoretical calculations have been performed for Ti-N (i.e., N_s
11 according to the notation used with the calculations with the model systems), Ti-NO
12 (i.e. N_i) and Ti-ON (i.e. N_{si}) species in *section 3.3*.

23 XPS has been also used to quantitatively estimate the concentration of nitrogen in the
24 films as a function of the plasma deposition conditions. Figure 4 shows a plot of the
25 N/Ti ratio as a function of the relative concentration of nitrogen in the plasma gas
26 determined by XPS for an ample series of thin films prepared at 523 K. In this plot we
27 make no distinction according to the precursor used for their synthesis, as the plasma
28 gas composition is the main parameter controlling the nitrogen content. It is also
29 important to mention that at 298 K the amount of incorporated nitrogen was always
30 slightly higher than at 523 K for N₂/O₂ ratios higher than 50% (data not shown). The
31 plot in Fig. 4 shows that the concentration of nitrogen in the film remains relatively
32 constant up to N₂ percentages of 80-85%. Above these values, the concentration of
33 nitrogen in the films sharply increases, reaching N/Ti ratios of ca. 0.5 when a 97%
34 nitrogen plasma was used for the synthesis. As mentioned previously, for this sample
35 the shape of the Ti2p spectrum broadens towards lower BEs, pointing to certain
36 reduction of the titanium cations and the formation of oxygen vacancies within the
37 lattice. The incorporation of nitrogen, particularly in the form of Ti-N species, and its
38 relationship with the formation/elimination of oxide vacancies within the lattice has
39 been addressed previously in the literature both experimentally and theoretically. ^{25, 34,}
40 ^{54, 55} In relation with these previous works dealing with an already formed TiO₂ network
41 where nitrogen is added either by ion implantation or upon annealing in ammonia, our
42 experimental conditions are different because the N-TiO₂ materials are created directly
43 by the deposition and/or reaction of/with plasma species. However, the fact that the
44 maximum concentration of Ti-N species appears when the oxide is defective suggests
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3 that the incorporation of this nitrogen species becomes favored in partially reduced
4 titanium oxide networks.
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7 8 *3.2 Optical and electrical properties of N-doped titanium oxide films* 9

10 The progressive incorporation of nitrogen into the films, particularly in the form of Ti-N
11 species, induces significant changes in their optical behavior. Fig. 5 shows a series of
12 UV-vis absorption spectra recorded for selected thin films prepared with TDMAT as
13 precursor and 85%, 90%, 95% and 97% N₂ in the plasma gas. It is important to note that
14 the three latter samples, have incorporated a considerable amount of nitrogen (cf. Figs. 3
15 and 4). These spectra are compared with the spectrum of a film prepared with 0% N₂
16 taken as a reference. All the curves are characterized by the typical oscillatory behavior
17 found when a high refraction index (**n**) material is deposited on a transparent substrate
18 with smaller **n**. These oscillations are originated by the interference of the partially
19 reflected beams at the film/air and film/substrate interfaces.⁵⁸ Besides these oscillations,
20 another quite remarkable effect is the shift of the absorption threshold to lower
21 wavelengths for thin films prepared with percentages of nitrogen higher than 85%.
22 Usually, this shift has been recognized in the literature as a hint of visible photo-activity
23 in doped TiO₂,¹⁷⁻²⁰ although this correlation has been claimed as a necessary but not
24 sufficient condition.^{22-24, 59} It is also important to remark in this figure that the minima
25 of the oscillations recorded for the sample prepared with 97% N₂ appear quite above the
26 curve of the reference TiO₂, which indicates a significant absorption in the visible
27 region (i.e., for $\lambda > 500$ nm). For percentages of N₂ in the plasma smaller than 85%, for
28 which the amount of incorporated nitrogen is still small (c.f. Fig.s 3 and 4), the shape of
29 the curves, except for the oscillatory behavior in the transparent region that depends on
30 the film thickness, was similar to that corresponding to the film prepared with 0% N₂
31 and did not depict any clear hint of a shifted absorption edge. Similar sets of curves
32 were obtained when using either TTIP or TDEAT as titanium precursor. For the whole
33 set of investigated samples, we have calculated the value of the absorption threshold by
34 assuming a Tauc model for the evaluation of the band gap of TiO₂ (i.e., $A(h\nu)^{1/2}$).⁶⁰ The
35 obtained values are represented in Fig. 6. The shape of the curve defined by the
36 different points is the reverse to that in Fig. 5, corresponding to the relative
37 concentration of nitrogen incorporated into the films. Consequently, a representation of
38 the absorption threshold values against the N/Ti ratio in the films follows a linear
39 tendency (see supporting information S3). This behavior indicates that the absorption
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3 threshold of the N-TiO₂ films is controlled by the amount and perhaps the type of
4 incorporated nitrogen. For the samples prepared with N₂ percentages higher than 85%
5 the Ti-N species, incorporated then in a substantial amount, is the most abundant
6 species, as evidenced by the shape of the XPS spectra in Figs. 2 and 3. These samples
7 present a clear shift in the absorption threshold whose magnitude is correlated with the
8 amount of nitrogen within the lattice (cf. Figs. 5 and 6). By contrast, samples prepared
9 with 85% or a lesser amount of nitrogen in the plasma, incorporate a limited amount of
10 nitrogen in the form of Ti-NO species. These samples do not present a clear shift in the
11 absorption threshold, although a clear determination is difficult with absorption spectra
12 such as those in Fig. 5, where the presence of interference oscillations may mask
13 changes at the edge spectral region.
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24 Electrical characterization of the films prepared with a high concentration of N₂ in the
25 plasma gas showed that they present a certain sheet conductivity (see supported
26 information S4) which is indicative of some reduction of the titanium ions and the
27 presence of a certain concentration of oxygen vacancies within the lattice.⁶¹ The broad
28 absorption feature extending through all the visible region of the spectra (i.e., $\lambda > 500$
29 nm) found for the samples prepared with 97% N₂ can therefore be associated with the
30 lost of stoichiometry in these samples.²¹ This confirms our evidences based on the
31 analysis of the Ti2p XPS spectra and supports the already commented relationship
32 between the incorporation of Ti-N species and the formation of oxygen vacancies within
33 the lattice.^{25, 34, 54, 55}
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43 Owing to the oscillatory behavior of the UV-visible transmission spectra, it is difficult
44 to ascertain if real absorption features close to the absorption threshold superimpose in
45 the spectra of the films prepared with percentages of N₂ lower or equal than 85%, where
46 only Ti-NO species have been determined by XPS. In order to check this point, we
47 employed ellipsometry to determine the extinction coefficient (**k**) and refraction index
48 (**n**) functions of selected films for $\lambda > 300$ nm. The corresponding plots are shown in
49 Fig. 7 for three thin films prepared at 523 K with TDMAT as precursor and percentages
50 of N₂ in the plasma gas of 0%, 85% and 97%. For the whole spectral region, the
51 refraction indices of the 97% N₂ film was higher than those of the other two samples
52 (e.g., at $\lambda = 550$ nm, the measured values were 2.38 and 2.12, respectively), in agreement
53 with its UV-vis absorption spectrum (Fig. 5) and its extinction coefficient curve plotted
54 in the bottom panel of Fig.6. This latter curve is characterized by a broad and intense
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3 maximum at around 400 nm that defines a shifted absorption threshold of 2.11 eV,
4 extending to cover the whole visible region of the spectrum. By contrast, the extinction
5 coefficient curve of the film prepared with 85% nitrogen does not present this broad
6 absorption and its edge jump defines an absorption threshold of 3.19 eV, quite similar to
7 that of the film prepared with pure oxygen as plasma gas. The main difference between
8 the curves of the samples prepared with 0% and 85% N₂ consists of a small but well
9 defined absorption feature at around 375 nm in the extinction coefficient curve of the
10 latter (indicated by a star in the figure). An absorption feature in this position was
11 indeed needed to extract a reasonable extinction coefficient curve from the ellipsometric
12 data recorded for this film.
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22 ***3.3. Density of states calculations for model N-TiO₂ systems***

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24 The three model systems used to calculate the density of states of the different N-TiO₂
25 thin films are shown in Fig. 8. They represent the substitutional, N_s , interstitial, N_i and
26 substitutional-interstitial, N_{si} defect states. In this latter type of defect state, N occupies
27 the same position as in the case of the N_s defect, but there is also an O atom occupying
28 an interstitial site. We will assume that these states of nitrogen correspond to the XPS
29 bands attributed in section 3.1, namely Ti-N, Ti-NO and Ti-ON species. Calculations
30 based on structures similar to those of the two former types of defects have been
31 previously carried out by Di Valentin et al.^{12,13, 15} to account for modifications of the
32 band structure in N-doped TiO₂. These authors have also considered either the effect of
33 Ti-N^{12,13, 15} or Ti-NO^{12, 13} like species, although they employed smaller simulation cells
34 (with 96 and 72 atoms for the anatase and rutile supercells respectively). Herein, besides
35 studying the effect of another type of species (i.e., Ti-ON), we address systematically
36 the possible influence in the electronic structure of the concentration of incorporated
37 nitrogen for both the Ti-N and Ti-NO species in a larger supercell, with 216 atoms. To
38 simulate the influence of the concentration of nitrogen in real samples, we have
39 considered two different simulation cells for each system, one containing one defect,
40 and another containing two defects. In addition, in this latter case we have used two
41 different arrangements for the two defect states, placing them in either two nearby or in
42 quite separated positions within the lattice. These two geometries are employed to
43 represent the electronic structure of N-TiO₂ systems with, respectively, a high and a low
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3 concentration of nitrogen. The most relevant geometrical parameters of the studied
4 defect configurations are reported in Table 1.
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8 We also performed calculations on a different type of defect, in which the O
9 atom was initially placed at its equilibrium lattice location, and the N atom is in an
10 interstitial site above the O atom at a distance of 1.35 Å along the z axis. We studied
11 two systems with one and two defects of this kind. Upon the energy minimization, the
12 cells underwent large geometrical rearrangements leading in the two cases to final
13 configurations equivalent to that of the N_i defects. This suggests that this initial
14 configuration is not stable and that minimization of the energy is accomplished by a
15 displacement of the O atoms from their initial location and their substitution by
16 nitrogen. This also means that the assumed detection by XPS of these species in some
17 of our samples (i.e., band at 401.7 eV of BE, cf. Fig. 3) is likely related with the fact
18 that the plasma deposition process utilized for the synthesis of the N-TiO₂ films is an
19 out of equilibrium procedure.
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30 The calculations also provided interesting information regarding the energy of
31 the different studied systems. Thus, when TiO₂ is modified with N_s defects, the energy
32 of the system where the two defect states are in nearby positions is 29.1 kJ/mol higher
33 than the energy of the system in which the two defects are far away. When considering
34 N_i defects, this energy difference is much smaller, just -1.1 kJ/mol, which means that
35 there is no strong energetic driving force preventing the aggregation of N_i defects. In the
36 case of N_{Si} defects, this energy difference is 14.0 kJ/mol, indicating that it is
37 energetically favourable to keep N_{Si} defects as far away as possible within the lattice.
38 On the other hand, to determine the relative stability of the N_{Si} and N_i defect states we
39 can compare directly the energies of the corresponding systems, as they have the same
40 number of atoms. For the systems modelling low nitrogen concentrations (with defects
41 far away from each other), we find that N_i defects are 8.3 kJ/mol more stable than N_{Si}
42 defects. Meanwhile, for the systems modelling a high concentration of nitrogen, the N_i
43 defects are 29.0 kJ/mol more stable than N_{Si} defects, a result which suggests that TiO₂
44 materials with a high concentration of nitrogen will tend to have N_s and N_i defects and
45 that the formation of N_{Si} defects would not be favoured.
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59 The total densities of states (DOS) calculations for pure TiO₂ (shown in the
60 Supporting Information S6) predict a gap between the valence and conduction bands of

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3 2.7 eV. This value is smaller than the experimental band gap of anatase TiO₂ (i.e., 3.2
4 eV⁶²), a difference that is due to the well known underestimation of the band gap in
5 DFT calculations.⁶³ The value of the Hubbard parameter *U* that we used is 4.5 eV,
6 which provides a good description of the gap states in periodic LDA+*U* calculations of
7 TiO₂ and, at the same time, predicts a reasonably good value for the band gap (i.e., 2.7
8 eV).⁴⁶ A similar result has been also obtained by Di Valentin et al.¹¹ when modeling N-
9 TiO₂ systems.

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17 The total densities of states (DOS) of the six N-TiO₂ studied systems are shown
18 in Fig. 9. The reported results indicate that the N-TiO₂ systems have developed new
19 electronic states in the band gap close to the valence states of titanium dioxide. A
20 similar result has been obtained by other authors simulating the electronic structure of
21 N-TiO₂ systems,^{11-15, 25, 34} although in these previous works smaller supercells were
22 employed in the calculations. Particularly interesting with regard to the optical
23 properties of N-TiO₂ is the difference observed in Fig. 9 between the shape of the gap
24 states corresponding to the *N_s* and the *N_i* and *N_{si}* defects around the HOMO level
25 (marked with a continuous line in the figure). In the former case (i.e. DOS curves a) and
26 b)) the extra states due to the incorporation of N appear as a continuous prolongation of
27 the valence band states of pure TiO₂ (marked with a dashed line in the figure). By
28 contrast for the case of *N_i* and *N_{si}* defect states, the N-related gap states appear as well
29 defined discrete peaks separated from the valence band of TiO₂ (i.e. DOS curves c)-f)).
30 Phenomenologically, this difference is equivalent to say that the *N_s* defects have
31 induced a shift of the valence band edge, producing a decrease in the gap of the system,
32 while the *N_s* and *N_{si}* states do not modify the valence band of TiO₂ but introduce new
33 electronic states in the gap. This point has been previously discussed based on both
34 experimental⁵⁴ and theoretical¹¹⁻¹⁵ studies. Thus, Diebold et al.⁵⁴ studied by XPS and
35 valence band photoemission N-implanted anatase and rutile, and concluded that
36 localized states rather than a band gap narrowing develop upon nitrogen implantation
37 and the subsequent formation of N³⁻ (i.e., Ti-N) species. Theoretically, studying the
38 anatase polymorph, Di Valentin et al.¹¹⁻¹² concluded that substitutional nitrogen states
39 (i.e., Ti-N) lie just above the valence band, while interstitial nitrogen states (i.e., Ti-NO)
40 lie higher in the gap. These authors also found that Ti-N species in rutile might produce
41 a certain increase of the gap with the N-associated states close to the valence band
42 maximum.¹⁵ The main difference between these previous studies and our calculations
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3 reside in the fact that we deal with more than one nitrogen species and define two
4 different scenarios for them depending on whether they are located in separated or close
5 positions of the lattice. When N atoms are close, as in the DOS curves a), c) and e), the
6 number of gap states is higher and appear more scattered along the band gap, as if there
7 is a splitting of peaks due to the electronic interaction between defects. By contrast,
8 when the N atoms are separated within the lattice, as in DOS curves b), d) and f), there
9 is a smaller number of peaks in the gap and they tend to be more concentrated around
10 certain energy values in the band gap (i.e., more localized states). In particular, for the
11 N_s species located in nearby lattice positions, it appears that the TiO_2 valence band
12 maximum extends continuously to higher energies. From an experimental point of view,
13 this is equivalent to say that there has been a narrowing of the gap. This would be in fact
14 the situation for our samples prepared with nitrogen percentages in the plasma higher
15 than 85% (cf. Figs. 4 and 5).
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27 As shown in the partial density of state plots in Fig. 10, the new states developed
28 in the gap have a mixed titanium-nitrogen-oxygen character and show slightly different
29 features depending on the kind of nitrogen defect considered in each case. This figure
30 represents in an enlarged scale the gap zone to clearly depict the new states appearing in
31 the N- TiO_2 systems. The plots represent the titanium, oxygen and nitrogen projected
32 states of the DOS curves in the regions around the HOMO lines in each system (note
33 that the HOMO lines appear at different energies in each case). For oxygen, they are
34 represented both the states corresponding to the oxygen atoms close to nitrogen and
35 those located in separated positions. This distinction is important since for the N_i and
36 N_{si} states nitrogen is directly bonded to an oxygen atom whose contribution to the DOS
37 of the new states is expected to be quite substantial. The first conclusion we could draw
38 from these plots is that all the defect states have a mixed character with contributions
39 from the oxygen, nitrogen and titanium atoms. However, a closer look to these curves
40 also reveals that in the N_s defects the nitrogen contribution to the gap states is relatively
41 higher than that of the close oxygen atoms, while for the N_i and N_{si} defects the nitrogen
42 and close-oxygen contributions are rather similar, thus supporting the importance of
43 these oxygen atoms in determining the electronic properties of the system. This closer
44 look to the new states around the HOMO line also suggests that the DOS curves of the
45 N_s defects, particularly for the case with the two nitrogen atoms located in nearby
46 positions, are a continuous extension of the valence band. The fact that the nitrogen
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3 contribution to these states extends through the entire DOS region, both below and
4 above the HOMO line, supports our view in the sense that a narrowing of the valence
5 band has occurred in this case. By contrast, for the N_i and N_{si} defects, the nitrogen
6 contribution to the new states appears quite concentrated in localized levels that do not
7 form a continuum with the valence band states. This result suggests that even for N-
8 TiO_2 systems with a high concentration of Ti-NO species, their electronic structure can
9 be described by the appearance of localized states in the gap.

16 3.4 Wetting angle photo-activity of N-TiO₂ thin films

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19 The photo-activity of N-TiO₂ thin films has been the subject of a large debate during the
20 last years.¹⁻¹⁴ In recent publications we have shown that transformation of the surface
21 state of N-TiO₂ from partially hydrophobic (i.e., wetting contact angles around 80°) into
22 hydrophilic can be induced by visible light in samples containing Ti-NO species,
23 regardless of whether they were amorphous or crystalline. By contrast, visible photo-
24 catalytic activity towards, for example, the photo-degradation of dyes seems to
25 constitute a more demanding process, which requires titanium oxide to be in crystalline
26 phase, forming a mixture composed mainly by anatase and a small amount of rutile.²²

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28 In agreement with previous studies with PECVD TiO₂ thin films using TTIP as volatile
29 precursor of titanium and pure oxygen as plasma gas²⁵), we have found that amorphous
30 TiO₂ is obtained for deposition temperatures lower than 523 K and that the anatase
31 phase of this oxide is obtained for $T \geq 523$ K. By contrast, most N-TiO₂ thin films
32 prepared at this temperature with TDMAT and TDEAT as precursors were amorphous,
33 as shown by their XRD spectra (see supporting information S7). Since no photo-
34 catalytic activity was found for any of the N-TiO₂ thin films when illuminated with
35 visible light and very little when using UV light, to check the possible role of the
36 different nitrogen species in controlling the visible photo-activity of these materials we
37 have looked for changes in wetting contact angle under visible and UV light
38 illumination. Fig. 11 presents some selected results showing that for samples containing
39 Ti-NO species (i.e., a sample prepared with 85% N₂) the wetting angle decreases by
40 almost 50° when illuminated with visible light, while for the films with a high
41 concentration of Ti-N species (i.e. sample prepared with 97% N₂) no visible light effect
42 was observed. This different behavior according to the type of nitrogen species was
43 common for the rest of N-TiO₂ thin films. When after the visible illumination the
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3 samples were irradiated with UV light, we found that all of them became
4 superhydrophilic (i.e., wetting angle lower than 10°), which is the typical behavior of
5 UV-illuminated TiO_2 thin films. Even if wetting induced changes and photo-catalytic
6 processes are not equivalent tests of TiO_2 photo-activity,^{22, 64} the reported results in Fig.
7 11 clearly prove that Ti-N species are unable to induce visible light photo-activity and
8 that the presence of Ti-NO species in the films is a requisite for promoting visible light
9 activation of N- TiO_2 thin films, at least for a partial hydrophilic conversion of their
10 surface.
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21 **4. Discussion**

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23 The previous results on N- TiO_2 thin films prepared by PECVD have provided
24 information about both the influence of the incorporated nitrogen on the optical
25 properties of TiO_2 and on the characteristics of the nitrogen species incorporated within
26 the structure of the oxide according to the different conditions of preparation. In relation
27 with this latter aspect, we have also described theoretically how the different species of
28 nitrogen may affect the electronic structure of the prepared materials. The influence of
29 these species of nitrogen on the photo-activity of N- TiO_2 is finally discussed in relation
30 with visible-light hydrophilic conversion of the surface of the prepared films.
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38 *4.1.-Type and concentration of nitrogen species and band gap narrowing in N- TiO_2 thin* 39 *films*

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42 The previous XPS results have shown that the incorporation of small amounts of
43 nitrogen in the form of Ti-NO species (i.e., N_i defects according to the notation used for
44 the density of states calculations) occurs when using TDMAT or TDEAT as precursors
45 or when the films are deposited with a O_2+N_2 plasma mixture with less than 90% of
46 nitrogen. We have proposed that the existence of four Ti-N bonds in the two latter
47 precursor molecules may account for the presence of Ti-NO species in these films, even
48 if no nitrogen gas has been added to the plasma. For 90% or higher concentration of
49 nitrogen a progressive incorporation of Ti-N species takes place. For all experimental
50 conditions used for the synthesis of the N- TiO_2 films we have also found that the
51 amount of incorporated nitrogen species (Ti-N and Ti-NO) was higher at low (i.e., 298
52 K) than at high (523 K) temperatures. The stability of Ti-N implanted in rutile and
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3 anatase polymorphs of TiO_2 has been addressed in previous articles,^{34, 54} where a
4 certain desorption upon annealing has been observed only for the case of anatase⁵⁴ and
5 an extra stabilization is achieved when gold is deposited onto the surface.³⁴ These
6 results have been interpreted in terms of the interaction of the implanted nitrogen (and
7 deposited gold particles) with oxygen vacancies in the lattice.^{25,34, 54} Although most of
8 our films were amorphous and a comparison with literature data is therefore not
9 straightforward, we must stress that our results point to a certain relationship between
10 the formation of Ti-N species and the existence of oxygen vacancies within the lattice.
11 The aforementioned dependence of the Ti-N concentration on the temperature of the
12 PECVD process and the fact that the maximum concentration of this species is found
13 for defective and high conductive samples support such dependence.
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24 Ti-N and Ti-NO species of nitrogen have been previously detected and discussed in
25 literature,^{1, 12, 22, 51-57} although no complete agreement exists on their specific role in
26 altering the electronic and optical properties of N- TiO_2 materials. In fact, during the last
27 years there has been a vivid discussion on whether modifying the TiO_2 with nitrogen
28 produces a shift in the absorption edge or just the development of absorption centers
29 close to an unmodified gap. Our experimental results indicate that both localized band
30 states and band gap narrowing may occur depending on the type and concentration of
31 nitrogen species incorporated within the films. Thus, we have found experimentally that
32 the incorporation of a relatively low concentration of interstitial nitrogen (i.e. Ti-NO
33 species) does not modify substantially the band gap structure of the oxide, but induces
34 the formation of some localized electronic states in the gap (cf. Figs. 5 and 7). A similar
35 situation has been claimed in the theoretical works of Di Valentin et al.¹¹⁻¹⁵ By contrast,
36 our data suggest that a narrowing of the band gap due to a shift in the valence band edge
37 may occur when substitutional nitrogen (i.e. Ti-N species) is incorporated within the
38 lattice. We have also found that the magnitude of this shift is directly related with the
39 amount of incorporated Ti-N species that we control by changing the relative amount of
40 nitrogen in the plasma. In recent experimental⁵⁴ and theoretical¹¹⁻¹⁵ studies on titanium
41 oxide with incorporated Ti-N species such a narrowing has not been detected, a
42 difference that might be related with the fact that in these experiments a lower
43 concentration of nitrogen species was actually incorporated within the lattice and
44 because only the effect of isolated nitrogen species was theoretically modeled.
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3 The films with a high concentration of nitrogen in the form of Ti-N species (i.e., those
4 prepared with a N₂ percentage of 95% or above) present a loss in transmission that has
5 been associated with the formation of oxygen vacancies (associated with Tiⁿ⁺ species,
6 n<4, detected by XPS) that would be responsible for a wide absorption extending over a
7 large zone of the visible spectrum.^{21, 61} The decrease in the resistivity of these films (cf.
8 supporting information S7) confirms that these films are oxygen defective, again
9 suggesting that incorporation of Ti-N species becomes favored by the presence of
10 oxygen vacancies.^{25, 34, 54}
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18 4.2. Density of states and electronic structure in N-TiO₂

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20 The density of states calculations reported in *section 3.4* account for the different
21 phenomenological behaviors found for N-TiO₂ thin films depending on the type and
22 concentration of the nitrogen species incorporated into the lattice. We have shown that
23 the N_i and N_{si} defects, associated to the Ti-NO and Ti-ON species detected by XPS, give
24 rise to localized electronic states in the gap above the valence band edge of the oxide. In
25 previous works, Di Valentin et al^{11,13, 15} reached a similar conclusion by simulating the
26 electronic structure of isolated N_i defect states. We have investigated the influence of
27 inter-defect distances on the electronic properties, in larger supercells. For a two-defect
28 model simulating the incorporation of a high concentration of N_i species within the
29 lattice (cf. Figs. 9 and 10), we have also proved that the gap states do not become
30 extensively mixed with the valence band levels of the oxide and keep a localized
31 character. This theoretical result supports the experimental findings by ellipsometry and
32 UV-vis spectroscopy for the 85% N₂ sample (cf. Fig. 7), where the presence of a
33 localized harmonic oscillator at around 375 nm is required for a good fitting of the
34 experimental curves. Unfortunately, our experimental protocol does not enable the
35 incorporation of a high concentration of Ti-NO species without Ti-N species being also
36 formed. This limitation might be explained by the calculated energies of formation of
37 the different defect states (*section 3.4*). Even though, the good agreement between
38 experimental and theoretical results supports the development of isolated states in the
39 gap rather than a band gap narrowing in the samples prepared with 85% or a lower
40 concentration of N₂ in the plasma gas.
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59 A direct comparison between our experimental and theoretical results is possible when
60 dealing with the N_s defect states (i.e., Ti-N species according to the notation used by the

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3 XPS analysis) where, particularly for the structures simulating a high concentration of
4 this species, a substantial mixing between the valence band states and those pertaining
5 to the nitrogen can be deduced from the calculated partial density of state curves (c.f.
6 Fig.10). As an overall effect, it can be interpreted that this mixing produces a shift
7 towards the visible of the valence band edge, in agreement with our experimental
8 observations for the samples prepared with nitrogen rich plasmas. As discussed in
9 *section 4.1*, the different interpretation regarding the modifications of the gap that was
10 outlined in previous theoretical¹¹⁻¹⁵ and experimental⁵⁴ works, might be connected with
11 the fact that they simulate or interpret experiments with more separated Ti-N species. In
12 this way, our description of the optical and electronic properties of N-TiO₂ systems with
13 a high concentration of Ti-N species would be in a better agreement with the original
14 suggestions of Asahi et al.¹ of a band gap narrowing than with those of these more
15 recent works.
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27 *4.3. Wetting photo-activity of N-doped TiO₂*

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29 Visible photo-activity of N-doped TiO₂ thin films is a very elusive effect with somehow
30 contradictory results in the literature, where the red-shift of the valence band edge has
31 been taken sometimes as a sufficient evidence warranting reactivity under visible
32 illumination.¹⁶ In recent works, it has been pointed out that although visible absorption
33 is a required condition, it is not sufficient to assure visible photo-activity of N-TiO₂.
34 Moreover, different conditions seem to be required for photo-catalysis and for light-
35 induced wetting angle changes over flat surfaces of this material.⁶⁴ The former is a
36 much more demanding process that, besides the presence of nitrogen species, requires
37 that the nitrogen modified titania is crystalline and complies with some additional
38 structural conditions, such as the simultaneous presence of both anatase and rutile
39 polymorphs in the photo-active material.²² Since most N-TiO₂ thin films prepared in the
40 present work were amorphous, we have investigated the effectiveness of either Ti-N or
41 Ti-NO species in inducing some kind of visible photo-activity by following the changes
42 in wetting angle upon illumination with visible light. The experiments carried out with
43 the N-TiO₂ thin films have shown that samples with a high concentration of Ti-N
44 species do not show visible photo-activity and keep their original wetting angle even
45 after a prolonged illumination with visible photons (cf. Fig. 11). By contrast, films
46 containing Ti-NO species do present some visible photo-activity, which was evidenced
47 by a partial and progressive decrease in the water contact angle of ca. 50° from an
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3 original wetting angle of 98° to a wetting angle of 52° after 60 min illumination. This
4 behavior, even if occurring with a slow rate, indicates that in the presence of this species
5 visible light is capable of photo-activating the surface of these titanium oxide samples.
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7 A thorough discussion of this kind of partial hydrophilic transformation of N-TiO₂
8 surfaces upon visible light irradiation can be found in reference 24.
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13 Our theoretical and experimental results have shown that a small concentration of Ti-
14 NO species in the films do not modify significantly the band gap of the oxide but
15 induces some localized electronic states close to the band edge. Visible light excitation
16 up to the conduction band of the electrons associated to these electronic states in the gap
17 should render a conduction band electron and a localized hole in the gap state. We
18 believe that precisely the localized character of this photo-hole is the factor precluding a
19 continuous operation of these films as a photo-catalyst as the mobility of these photo-
20 holes induced with visible light is rather limited.¹⁵ Indeed, for a continuous photo-
21 catalytic process to take place it is required the migration of both electrons and holes
22 towards the surface, as it usually happens in undoped TiO₂ illuminated with UV light.
23 In N-TiO₂ systems, migration towards the surface of the photo-holes generated in the
24 illuminated layer (i.e. in the order of the micron) can only take place if a continuous
25 band-like state forms in the gap. This is not certainly the case in our samples with a low
26 concentration of incorporated Ti-NO species. However, as previously discussed by us,²⁴
27 visible-light partial conversion of the surface hydrophilicity occurs by just activating the
28 utmost atomic layers whereby only surface electron-hole pairs are involved in the
29 transformation. Therefore, the point to be stressed here is that while Ti-NO species are
30 able to induce this visible-light transformation, Ti-N species aren't. This finding
31 contrasts with the original proposal of Asahi et al.¹ and thereafter of other authors^{5, 16}
32 attributing to the Ti-N species the visible photo-activity of N-TiO₂ materials. However,
33 it complies with more recent results stressing that for visible photo-activity to occur, Ti-
34 NO like species must be present.^{26, 56}
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52 53 **5.-Conclusions**

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55 The previous results have shown that by using the PECVD technique it is possible to
56 obtain a large set of N-TiO₂ thin films where the amount and type of nitrogen can be
57 tuned by controlling the working parameters during the deposition of the films. Thus,
58 Ti-NO species are the only type of incorporated nitrogen species when the films were
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3 prepared with a relatively low concentration of nitrogen in the plasma gas. On the other
4 hand, a high concentration of Ti-N species becomes incorporated into the films for
5 nitrogen rich plasmas.
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9 Density of states calculations of differently doped anatase model systems have provided
10 a detailed description of the influence of the N states, Ti-N and Ti-NO (and the related
11 one Ti-ON), on the electronic structure of N-TiO₂. It is found that in anatase with a
12 relatively high concentration of Ti-N species the valence band edge shifts upwards
13 because of the mixing of the N atomic levels with the levels of the lattice atoms.
14 Consequently, the absorption edge of the UV-vis absorption spectra is red shifted as
15 found experimentally. On the contrary, for the Ti-NO species the electronic levels of the
16 interstitial nitrogen species remain localized and no band gap narrowing occurs.
17 Experimentally, this situation yields spectra whose absorption edge are not shifted but
18 present a separated absorption oscillator close to it.
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28 Films containing Ti-N species are not photo-active with visible light even if they
29 present a red-shifted absorption edge. Visible light wetting photo-activity is only
30 observed in N-TiO₂ films with Ti-NO defect states. This observation has been related
31 with the surface excitation of the electronic states associated to these species, which
32 appear in the gap close to a virtually un-modified valence band edge.
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43 **6.-Acknowledgments.-** We thank the Ministry of Science and Education of Spain
44 (projects MAT 2007-65764 and the CONSOLIDER INGENIO 2010-CSD2008-00023)
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46 support. This work has been carried out within the EU project NATAMA (contract no.
47 032583)
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52 **Supporting Information Available.**

53 Some additional information referred in the text as S1-S8 is available free of charge via
54 the Internet at <http://pubs.acs.org>.
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Table 1.- Relevant distances (in Å) involved in the three types of defects studied: N_i , N_s and N_{si} . The distance between the two defects in the unit cell is taken into account by labeling with C and S the structures in which the N atoms are close and separated respectively.

Species	O-Ti	N-O	N-Ti
N_i -C	2.06, 2.06, 2.01	1.48	1.95, 1.98, 2.05
N_i -S	2.13, 2.13, 2.28	1.36	2.06, 2.06, 2.28
N_s -C			1.98, 2.00, 2.07
N_s -S			2.00, 2.00, 2.09
N_{si} -C	2.09, 2.17, 3.33	1.29	1.98, 2.02, 2.04
N_{si} -S	2.17, 2.17, 3.30	1.28	2.02, 2.03, 2.03

Figure captions.-

Fig. 1.- Schematic representation of the Ti precursors utilized for the synthesis of the N-TiO₂ thin films by plasma deposition.

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3 Fig. 2.- N1s spectra recorded for a series of N-TiO₂ thin films prepared with the
4 TDMAT precursor and different percentages of N₂ in the plasma as indicated. Note that
5 some spectra are affected by a multiplication factor to bring them into a common scale.
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11 Figure 3.- Fitted (gray line) and experimental (black line) N1s normalized spectra of
12 selected N-TiO₂ thin films prepared at 523 K with the TDMAT (left) and the TTIP
13 (right) precursors and different percentages of N₂ in the plasma gas. Elemental bands
14 used for fitting are plotted in gray.
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18 Figure 4.- N/Ti ratio determined from the intensities of the N1s and Ti2p photoelectron
19 peaks as a function of the percentage of nitrogen in the O₂+N₂ mixture used as plasma
20 gas. The dashed line is included to guide the eyes.
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24 Figure 5.- Set of absorption spectra recorded for a series of N-TiO₂ thin films prepared
25 at 523 K with the TDMAT precursor and different percentages of nitrogen in the plasma
26 gas as indicated in each particular panel. The gray color spectrum included for
27 comparison in each panel corresponds to the absorption spectrum of a pure TiO₂ thin
28 film.
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34 Figure 6.- Values of the absorption threshold plotted against the percentage of nitrogen
35 in the plasma gas for N-TiO₂ thin films prepared with the TTIP, TDMAT and TDMAT
36 precursors at 523 K. The dashed line is included to guide the eyes.
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40 Figure 7.-Top) Refraction index and bottom) extinction coefficient curves determined
41 for N-TiO₂ thin films prepared at 523 K by using TDMAT as a precursor and different
42 percentages of N₂ in the mixture used as plasma gas as indicated. The star points to the
43 region whose reproduction is only possible by admitting an oscillator around 375 nm.
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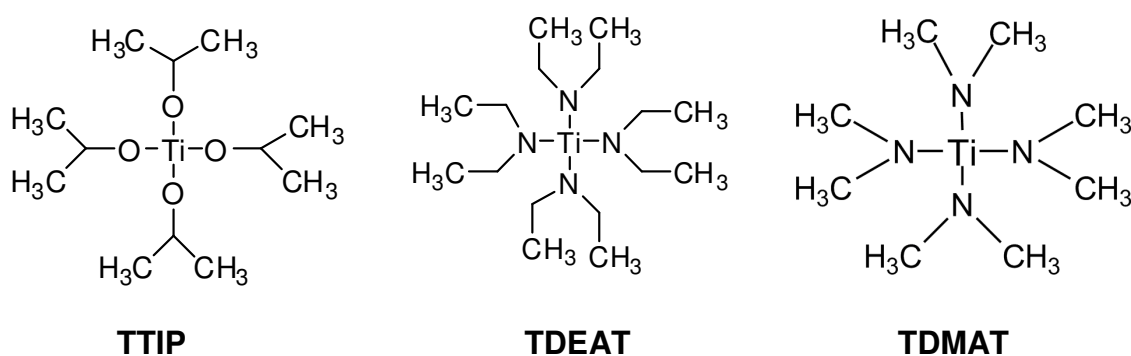
48 Figure 8.- Structures of the three type of N centers studied. a) Substitutional defect, N_s , in
49 which a N atom replaces a O atom. b) Interstitial defect, N_i , in which the N atom is located in
50 an interstitial place, and the O atom is also displaced from its original site. c) Substitutional-
51 interstitial defect, N_{st} , in which the N atom occupies the same position as in the case of the
52 N_s defect, and the O atom occupies an interstitial place.
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59 Figure 9.- Density of states plots of the six N-TiO₂ systems studied, showing the two
60 spin components separately. **a** and **b** correspond to the two systems with substitutional

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3 N doping (N_s), with 2 N atoms located at close and separated positions, respectively. **c**
4 and **d** correspond to the two systems with interstitial N doping (N_i), with 2 N atoms
5 located at close and separated positions, respectively, respectively. **e** and **f** correspond to
6 the two substitutional-interstitial N configurations (N_{si}), with 2 N atoms located at
7 close and separated positions, respectively, respectively. The bands are plotted in such a
8 way that the origin of the energy (0 eV) is placed at the bottom of the conduction band.
9 The two vertical dashed lines are placed to indicate HOMO-LUMO gap in TiO₂. The
10 vertical solid line indicate the highest occupied orbital of the system, which in the case
11 of N-TiO₂ corresponds to new gap states generated by the presence of nitrogen in the
12 lattice.
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23 Figure 10.- Partial density of states plots of N, Ti and O in the gap states for the six
24 studied N-TiO₂ systems with the labelling a)-f) being the same than in Fig. 9. The oxygen
25 contribution is differentiated for the oxygen atoms close to the nitrogen and the rest of
26 oxygen atoms of the lattice. The total densities of states plotted in Figure 9 are calculated
27 as the sum of these partial densities of states. The vertical solid line indicates the highest
28 occupied orbital of the system.
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34 Figure 11.- Evolution of the wetting angle as a function of the illumination time with
35 visible and UV lights and then left in the dark for N-TiO₂ samples prepared with 85%
36 (top) and 97% (bottom) of N₂ in the plasma gas. The lines are plotted to guide the eyes.
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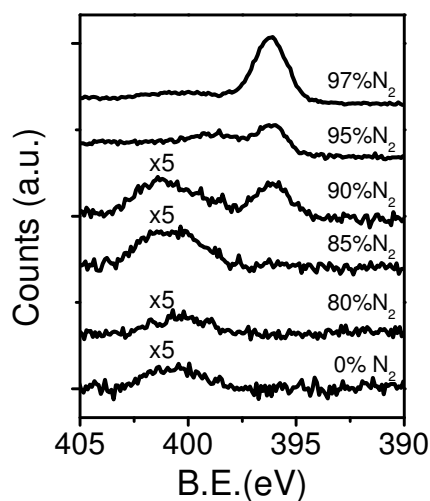


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18 **TTIP: titanium tetrakis isopropoxide**

19 **TDEAT: tetrakis diethyl amino titanium**

20 **TDMAT: tetrakis dimethyl amino titanium**

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27 Fig. 1.- Schematic representation of the Ti precursors utilized for the synthesis of the N-
28 TiO₂ thin films by plasma deposition.
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53 Fig. 2.- N1s spectra recorded for a series of N-TiO₂ thin films prepared with the
54 TDMAT precursor and different percentages of N₂ in the plasma as indicated. Note that
55 some spectra are affected by a multiplication factor to bring them into a common scale.
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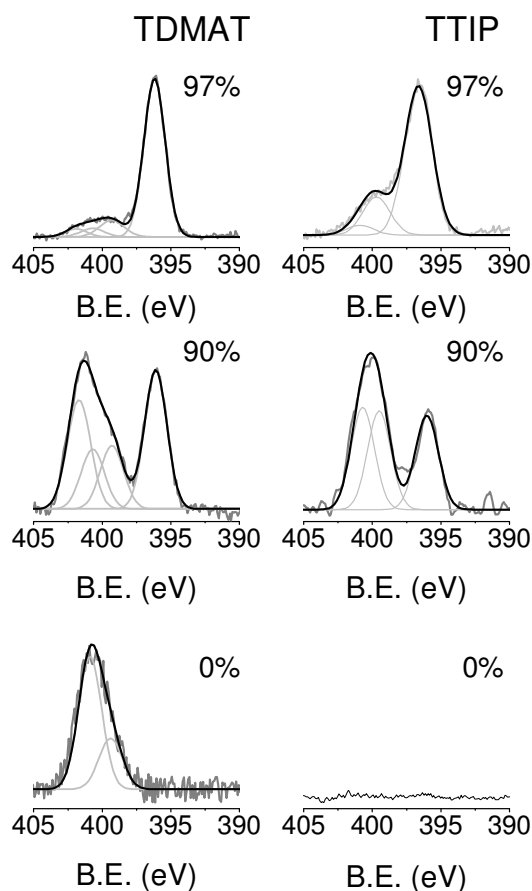


Figure 3.- Fitted (gray line) and experimental (black line) N1s normalized spectra of selected N-TiO₂ thin films prepared at 523 K with the TDMAT (left) and the TTIP (right) precursors and different percentages of N₂ in the plasma gas. Elemental bands used for fitting are plotted in gray.

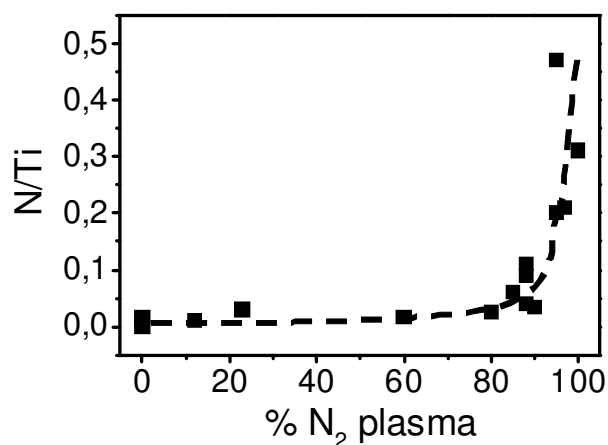


Figure 4.- N/Ti ratio determined from the intensities of the N1s and Ti2p photoelectron peaks as a function of the percentage of nitrogen in the O₂+N₂ mixture used as plasma gas. The dashed line is included to guide the eyes.

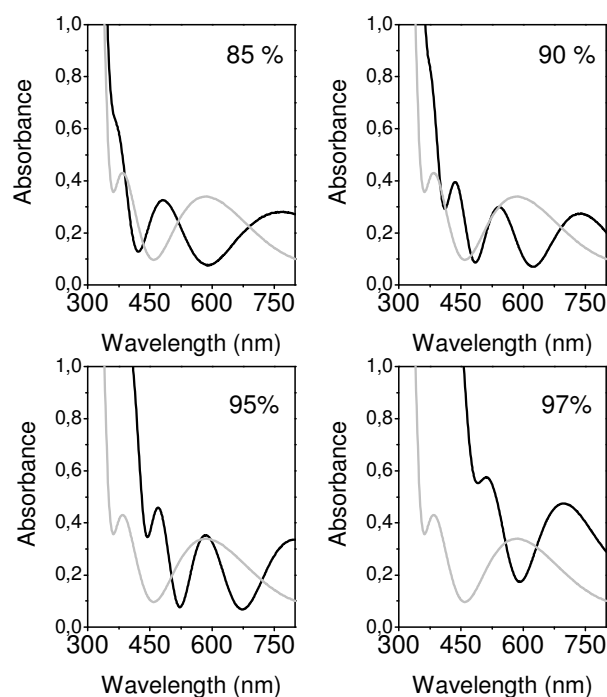


Figure 5.- Set of absorption spectra recorded for a series of N-TiO₂ thin films prepared at 523 K with the TDMAT precursor and different percentages of nitrogen in the plasma gas as indicated in each particular panel. The gray color spectrum included for comparison in each panel corresponds to the absorption spectrum of a pure TiO₂ thin film.

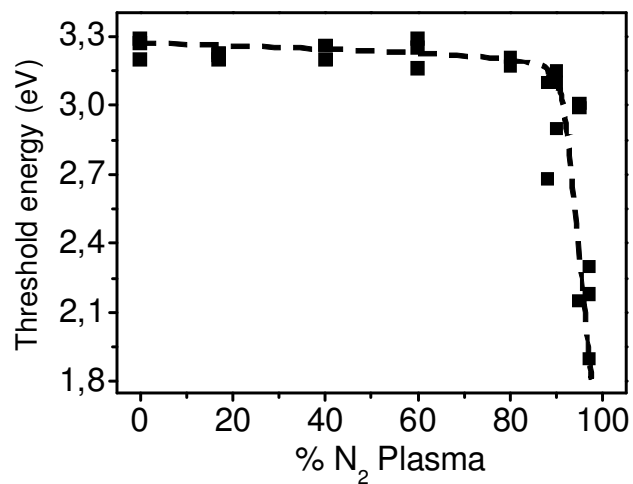


Figure 6.- Values of the absorption threshold plotted against the percentage of nitrogen in the plasma gas for N-TiO₂ thin films prepared with the TTIP, TDMAT and TDMAT precursors at 523 K. The dashed line is included to guide the eyes.

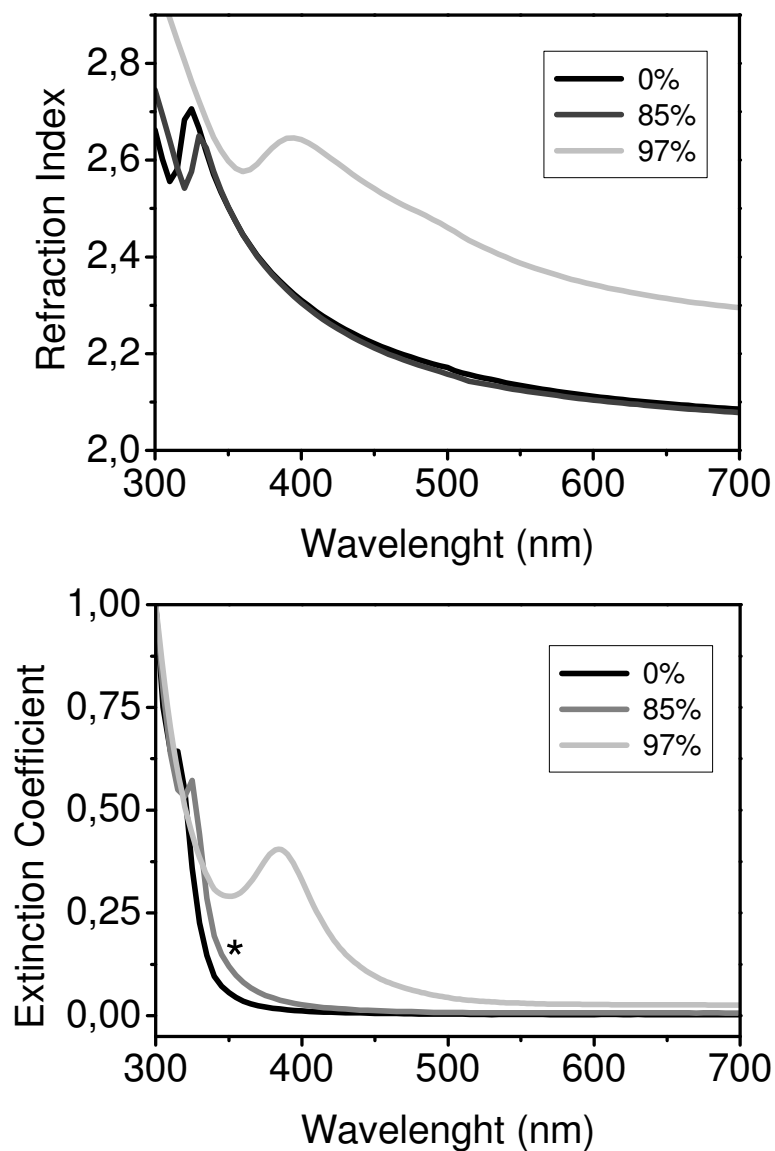


Figure 7.-Top) Refraction index and bottom) Extinction coefficient curves determined for N-TiO₂ thin films prepared at 523 K by using TDMAT as a precursor and different percentages of N₂ in the mixture used as plasma gas as indicated. The star points to the region whose reproduction is only possible by admitting an oscillator around 375 nm.

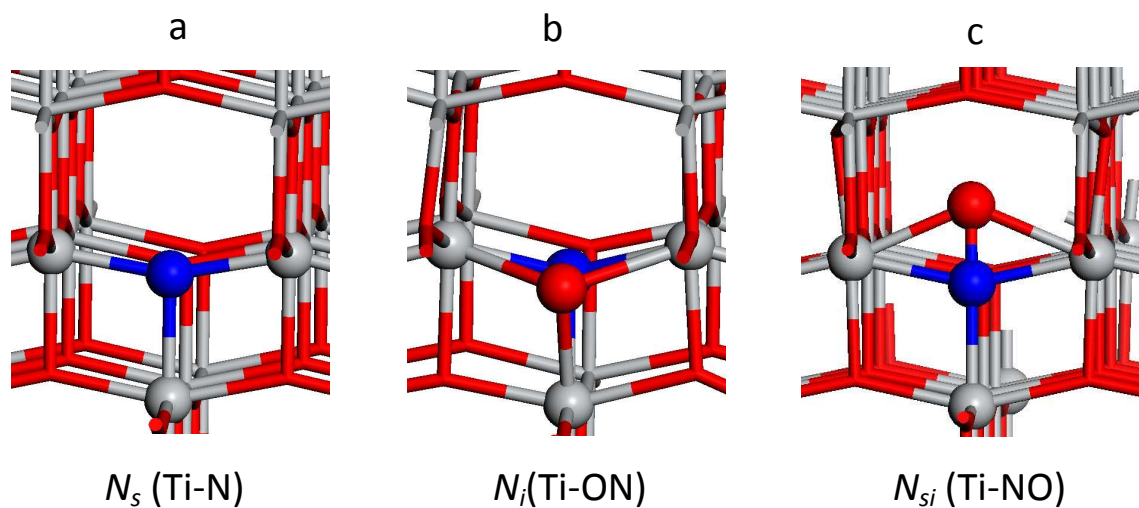


Figure 8.- Structures of the three type of N centers studied. a) Substitutional defect, N_s , in which a N atom replaces a O atom. b) Interstitial defect, N_i , in which the N atom is located in an interstitial place, and the O atom is also displaced from its original site. c) Substitutional-interstitial defect, N_{si} , in which the N atom occupies the same position as in the case of the N_s defect, and the O atom occupies an interstitial place.

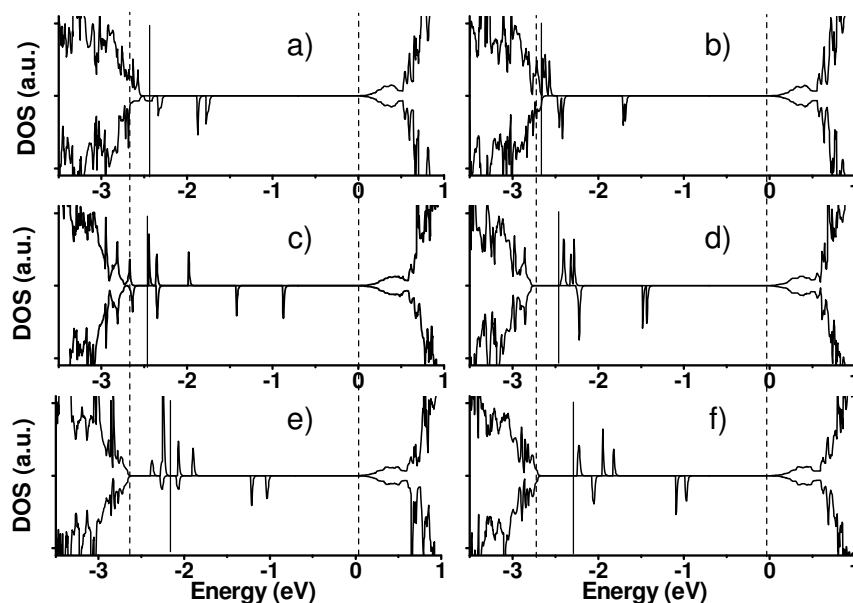


Figure 9.- Density of states plots of the six N-TiO₂ systems studied, showing the two spin components separately. **a** and **b** correspond to the two systems with substitutional N doping (N_s), with 2 N atoms located at close and separated positions, respectively. **c** and **d** correspond to the two systems with interstitial N doping (N_i), with 2 N atoms located at close and separated positions, respectively, respectively. **e** and **f** correspond to the two substitutional-interstitial N configurations (N_{si}), with 2 N atoms located at close and separated positions, respectively, respectively. The bands are plotted in such a way that the origin of the energy (0 eV) is placed at the bottom of the conduction band. The two vertical dashed lines are placed to indicate HOMO-LUMO gap in TiO₂. The vertical solid line indicate the highest occupied orbital of the system, which in the case of N-TiO₂ corresponds to new gap states generated by the presence of nitrogen in the lattice.

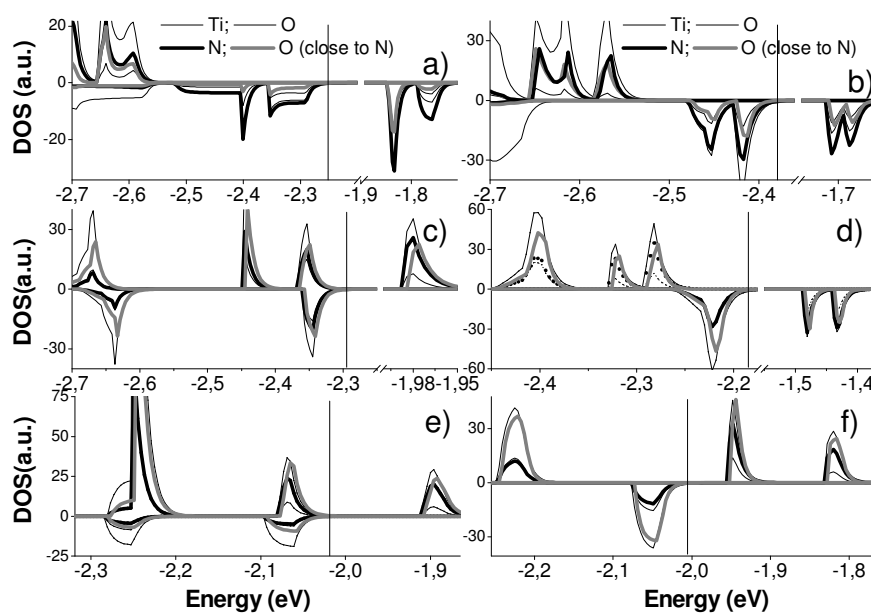


Figure 10.- Partial density of states plots of N, Ti and O in the gap states for the six studied N-TiO₂ systems with the labelling a)-f) being the same than in Fig. 9. The oxygen contribution is differentiated for the oxygen atoms close to the nitrogen and the rest of oxygen atoms of the lattice. The total densities of states plotted in Figure 9 are calculated as the sum of these partial densities of states. The vertical solid line indicates the highest occupied orbital of the system.

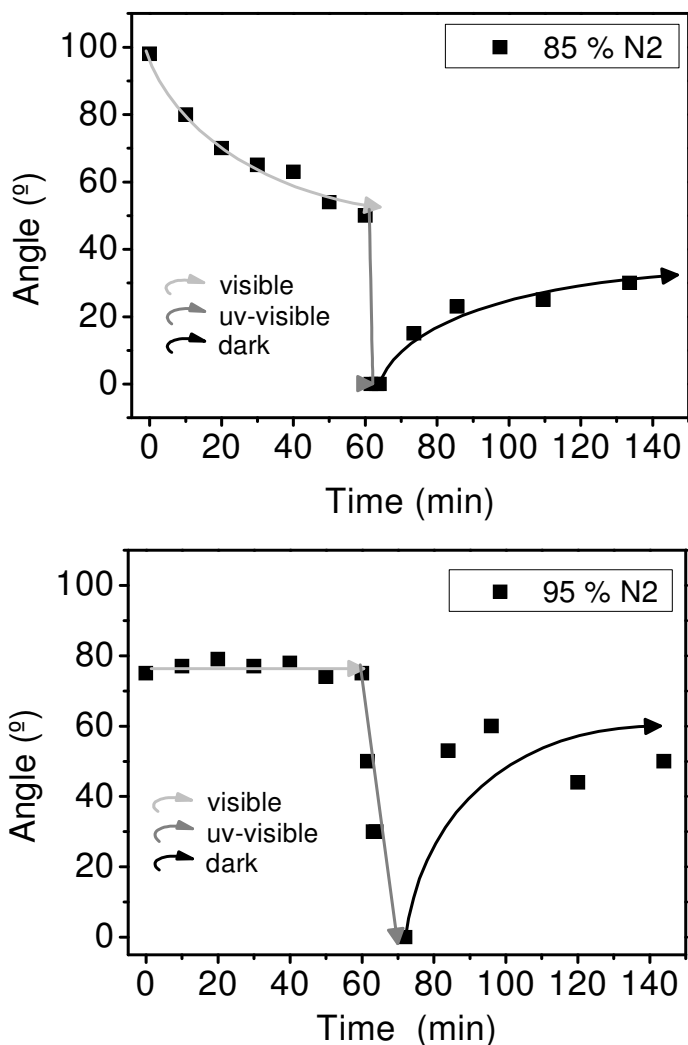
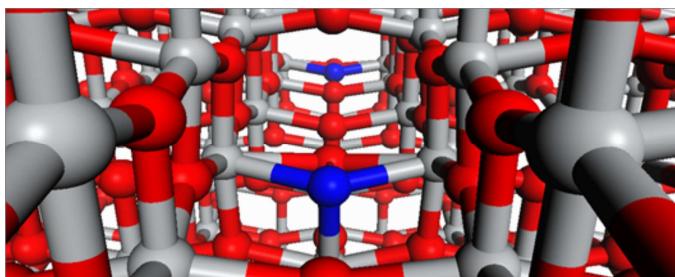


Figure 11.- Evolution of the wetting angle as a function of the illumination time with visible and UV lights and then left in the dark for N-TiO₂ samples prepared with 85% (top) and 97% (bottom) of N₂ in the plasma gas. The lines are plotted to guide the eyes.

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